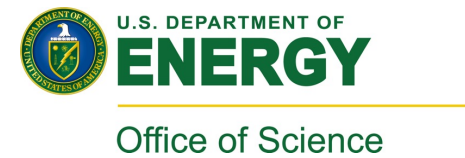


Evolution of cloud droplet spectral width: A new look at an old problem

Wojciech W. Grabowski

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This material is based upon work supported by the National Center for Atmospheric Research, which is a major facility sponsored by the National Science Foundation under Cooperative Agreement No. 1852977.

Adiabatic evolution of cloud droplet spectral width: A new look at an old problem

Wojciech W. Grabowski¹ and Hanna Pawlowska²

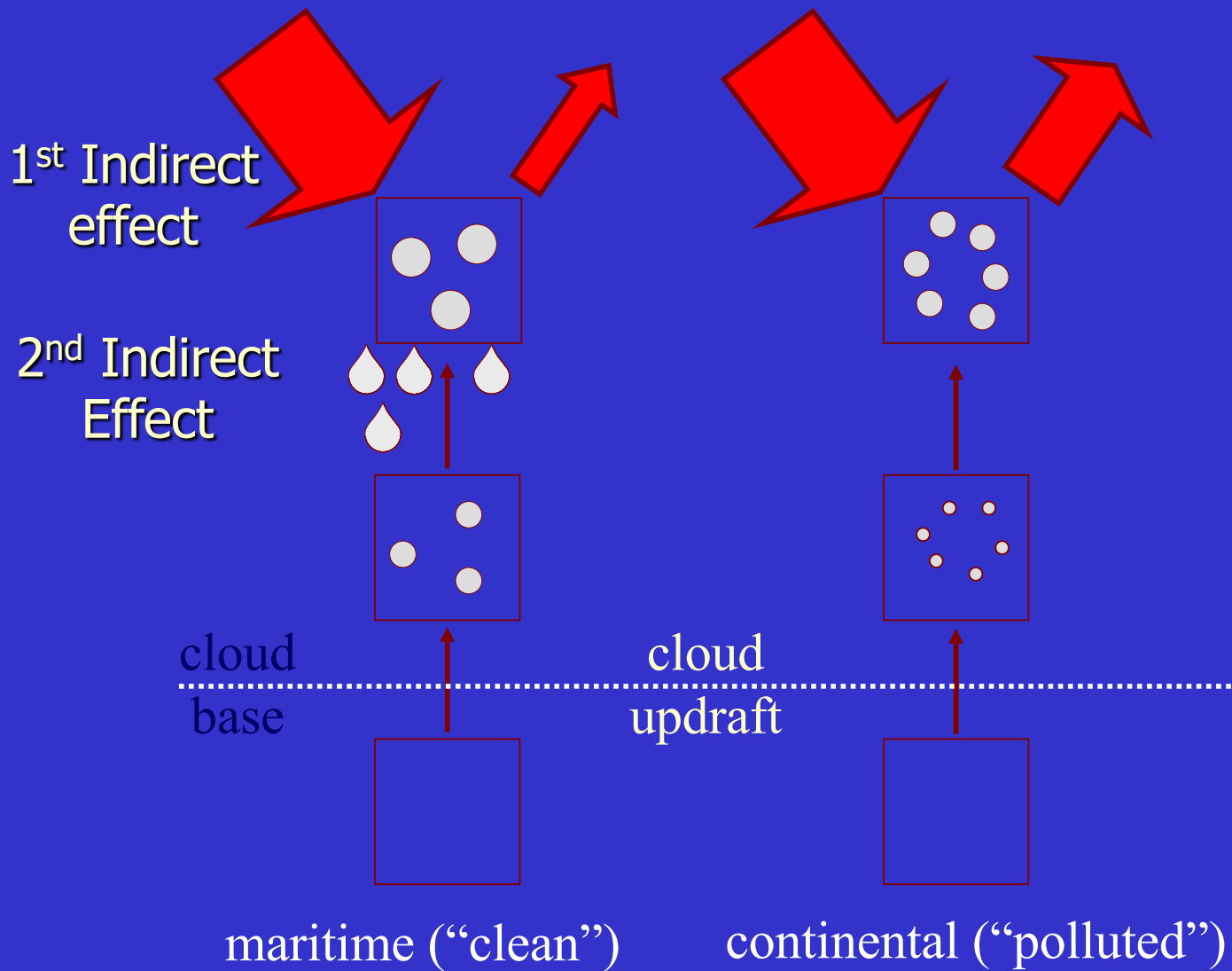
(manuscript in review for *Geophys. Res. Lett.*)

Why should one worry about width of the cloud droplet spectrum?

Impact on cloud radiative properties, especially for solar radiation.

Impact of rain development in ice-free clouds.

Indirect aerosol effects



Radiative properties (optical thickness, etc.) of warm clouds depend on the so-called **effective radius**, the ratio of the third and second moment of the droplet size distribution:

$$R_{eff} = \frac{\langle R^3 \rangle}{\langle R^2 \rangle}$$

Effective radius R_{eff} depends on the **mean volume radius** $\sim(\text{LWC}/N)^{1/3}$ and the droplet spectrum **relative dispersion** d , the ratio of the standard deviation of the droplet radius distribution and the mean droplet radius (Pontikis and Hicks GRL 1992; “PH” below; Liu and Daum GRL 2000):

effective radius

$$R_{eff} = \alpha \left(\frac{L}{N} \right)^{1/3}$$

L – LWC

N – droplet concentration

$d = \sigma / \langle R \rangle$

relative dispersion

$$\alpha_{PH}(d) = 62.04 \frac{(1 + 3d^2)^{2/3}}{(1 + d^2)}$$

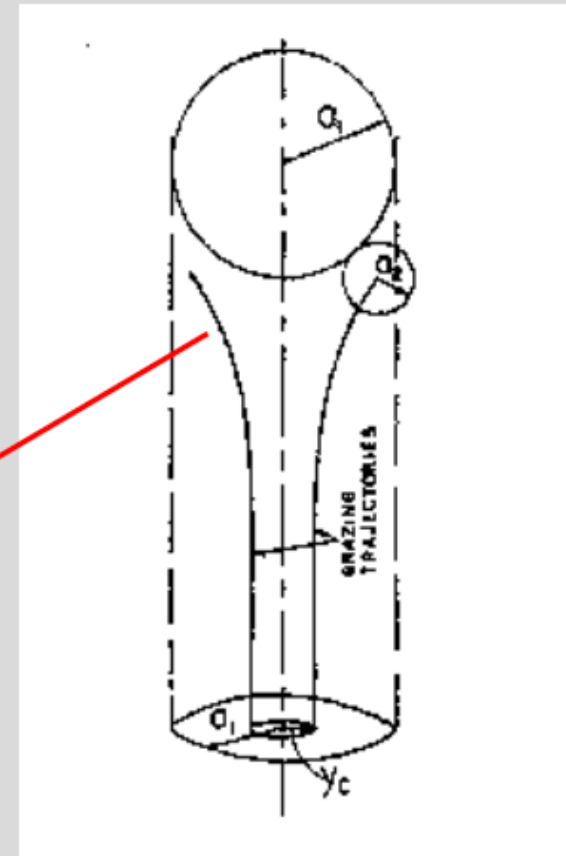
Growth of water droplets by gravitational collision-coalescence:

$$K(m_{a1}, m_{a2}) = E_c \pi (a_1 + a_2)^2 |(V_{a1} - V_{a2})|$$

Collision efficiency:

$$E_c = \frac{y_c^2}{(a_1 + a_2)^2}$$

Grazing trajectory



Droplet inertia is the key; without it, there will be no collisions. This is why collision efficiency for droplets smaller than 10 μm is very small.

Growth of water droplets by gravitational collision-coalescence:

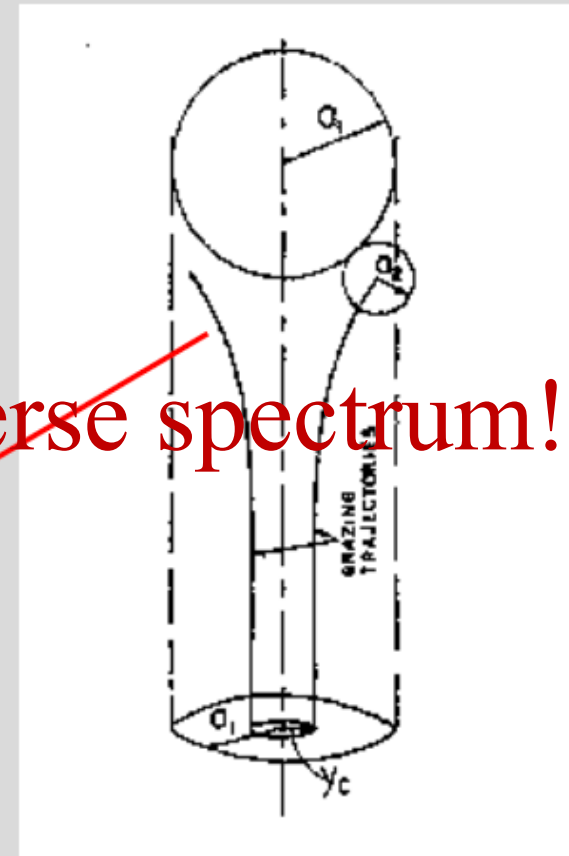
$$K(m_{a1}, m_{a2}) = E_c \pi(a_1 + a_2)^2 (|V_{a1} - V_{a2}|)$$

no collisions for a monodisperse spectrum!

Collision efficiency:

$$E_c = \frac{y_c^2}{(a_1 + a_2)^2}$$

Grazing trajectory



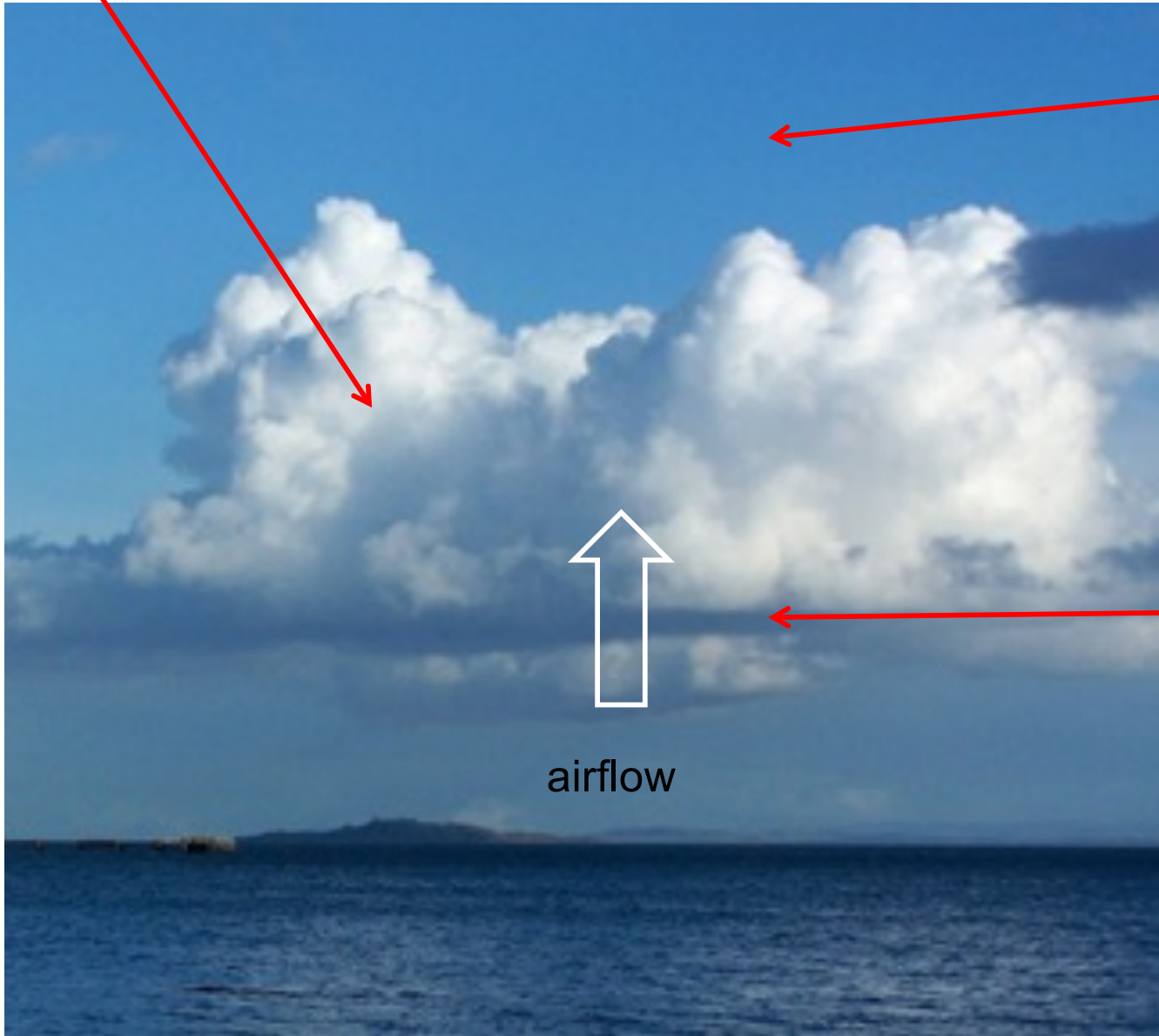
Droplet inertia is the key; without it, there will be no collisions. This is why collision efficiency for droplets smaller than 10 μm is very small.

turbulent
cloud

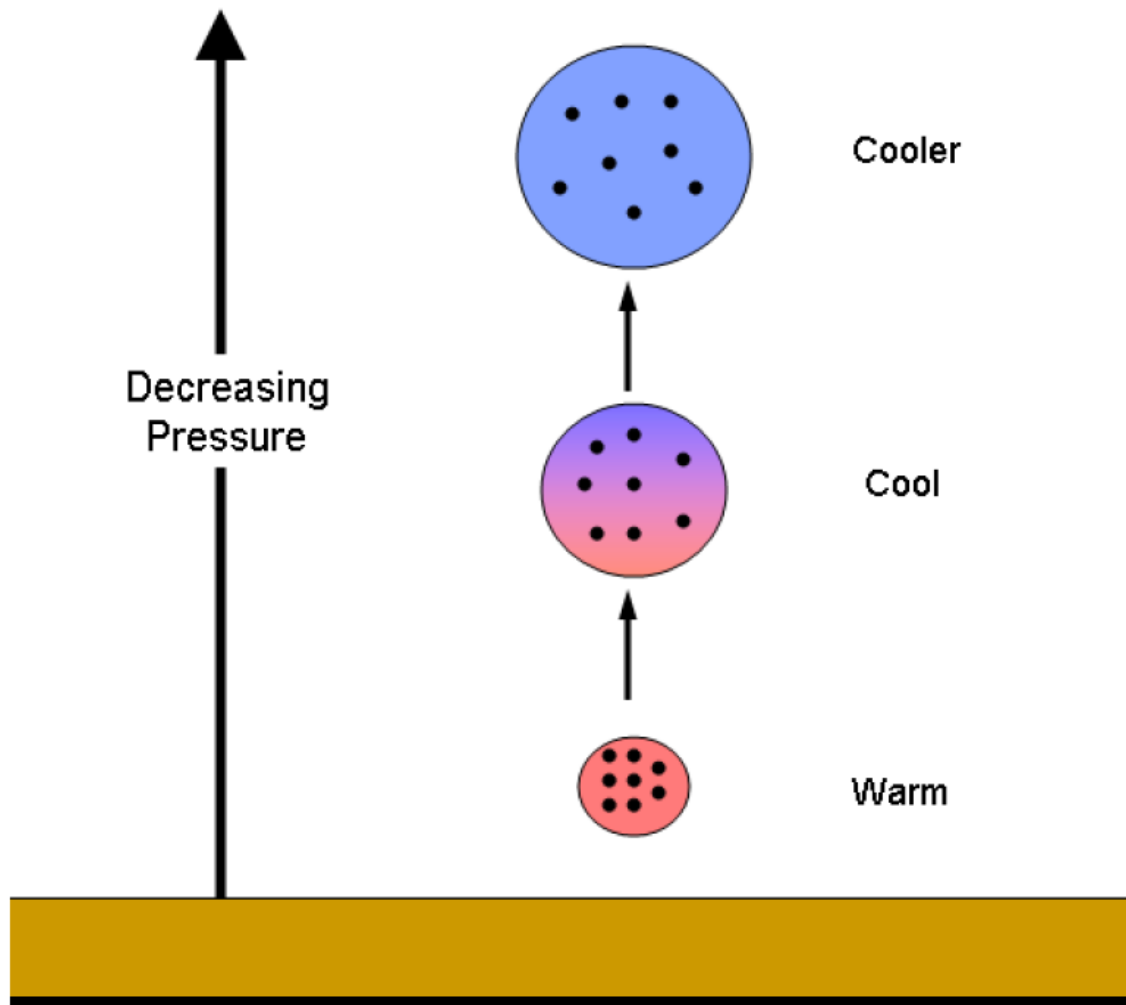
calm (low-
turbulence)
environment

cloud base:
activation of cloud
droplets

airflow



Traditional approach to study cloud-base activation:
adiabatic parcel crossing the cloud base with a prescribed ascent

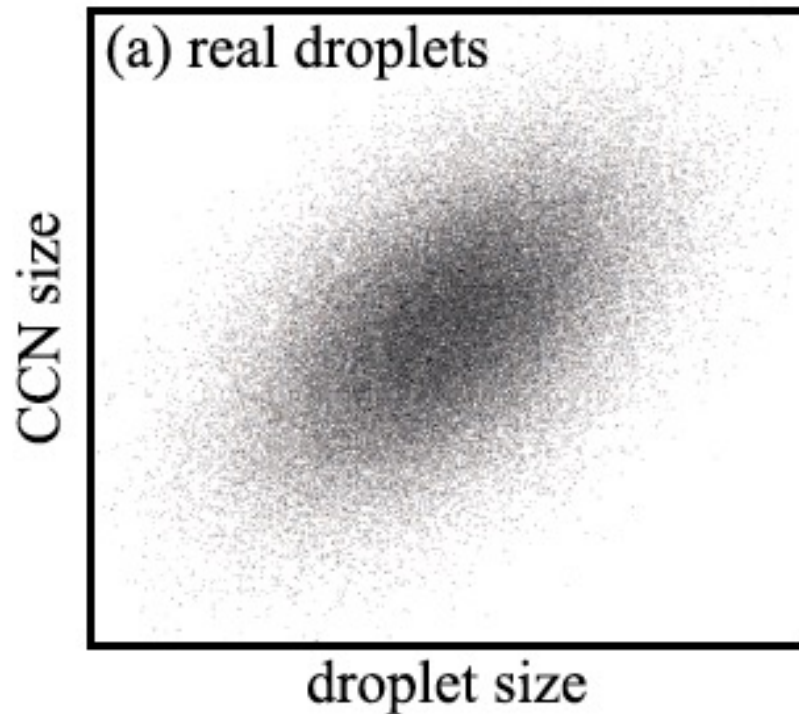


$$c_p \frac{dT}{dt} = -gw + L_v C_d$$

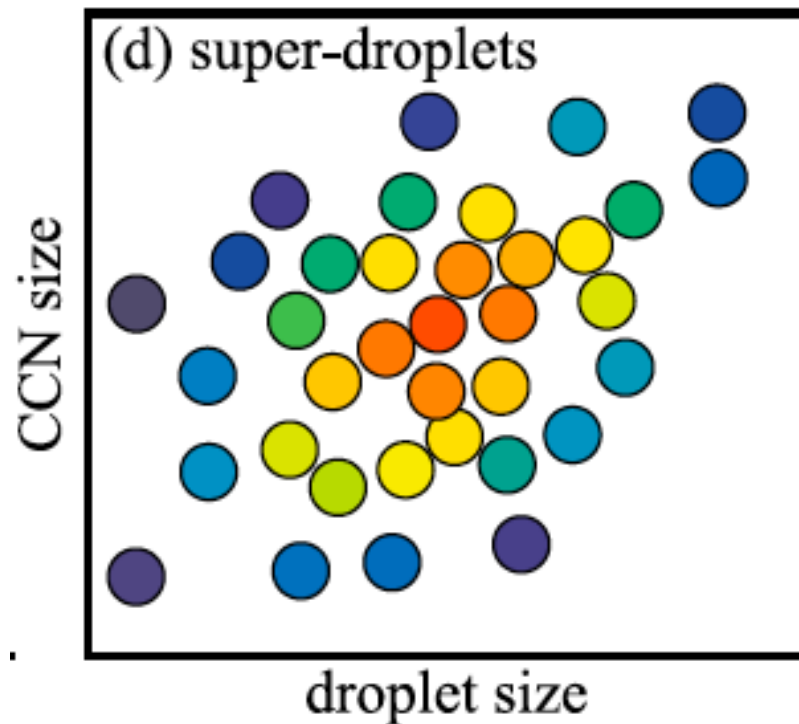
$$\frac{dq_v}{dt} = -C_d$$

$$\frac{dq_c}{dt} = C_d$$

$$\frac{dp}{dt} = -\rho w g$$

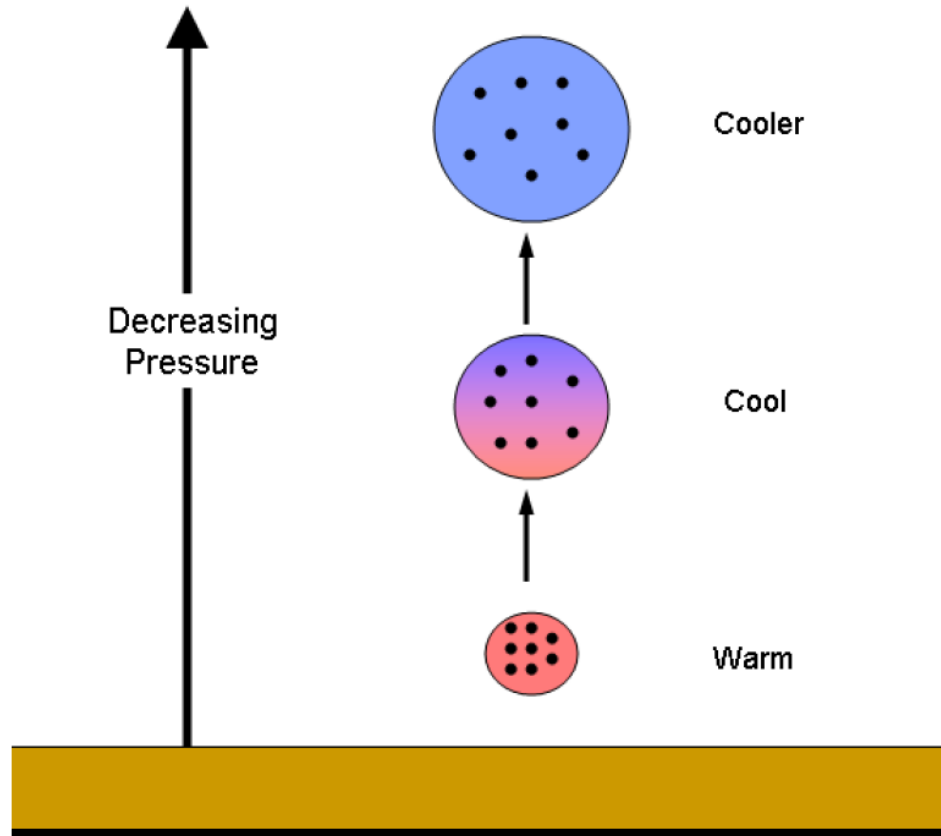


Lagrangian
microphysics:
super-droplets



Color represents
“multiplicity
attribute”
(red - high, blue – low)

Traditional approach to study cloud-base activation:
adiabatic parcel crossing the cloud base with prescribed ascent



$$c_p \frac{dT}{dt} = -gw + L_v C_d$$

$$\frac{dq_v}{dt} = -C_d$$

$$\frac{dq_c}{dt} = C_d$$

$$\frac{dp}{dt} = -\rho w g$$

Adiabatic parcel model has been used in numerous studies linking physicochemical CCN properties below cloud base to the cloud droplet concentrations aloft. Here we use the model to look at the evolution of the droplet spectral width.

adiabatic parcel crossing the cloud base:

Starting from $\text{RH} = 97\%$, $p = 900 \text{ hPa}$, $T = 283 \text{ K}$

CCN (superdroplets) from prescribed distribution,
assumed initially at equilibrium with ambient RH

prescribed ascent of $0.25, 1, \text{ or } 4 \text{ m s}^{-1}$

Deliquesced CCN / droplet growth equation:

$$d(r^2)/dt = G (S - S_{eq})$$

S - ambient supersaturation

S_{eq} - equilibrium supersaturation

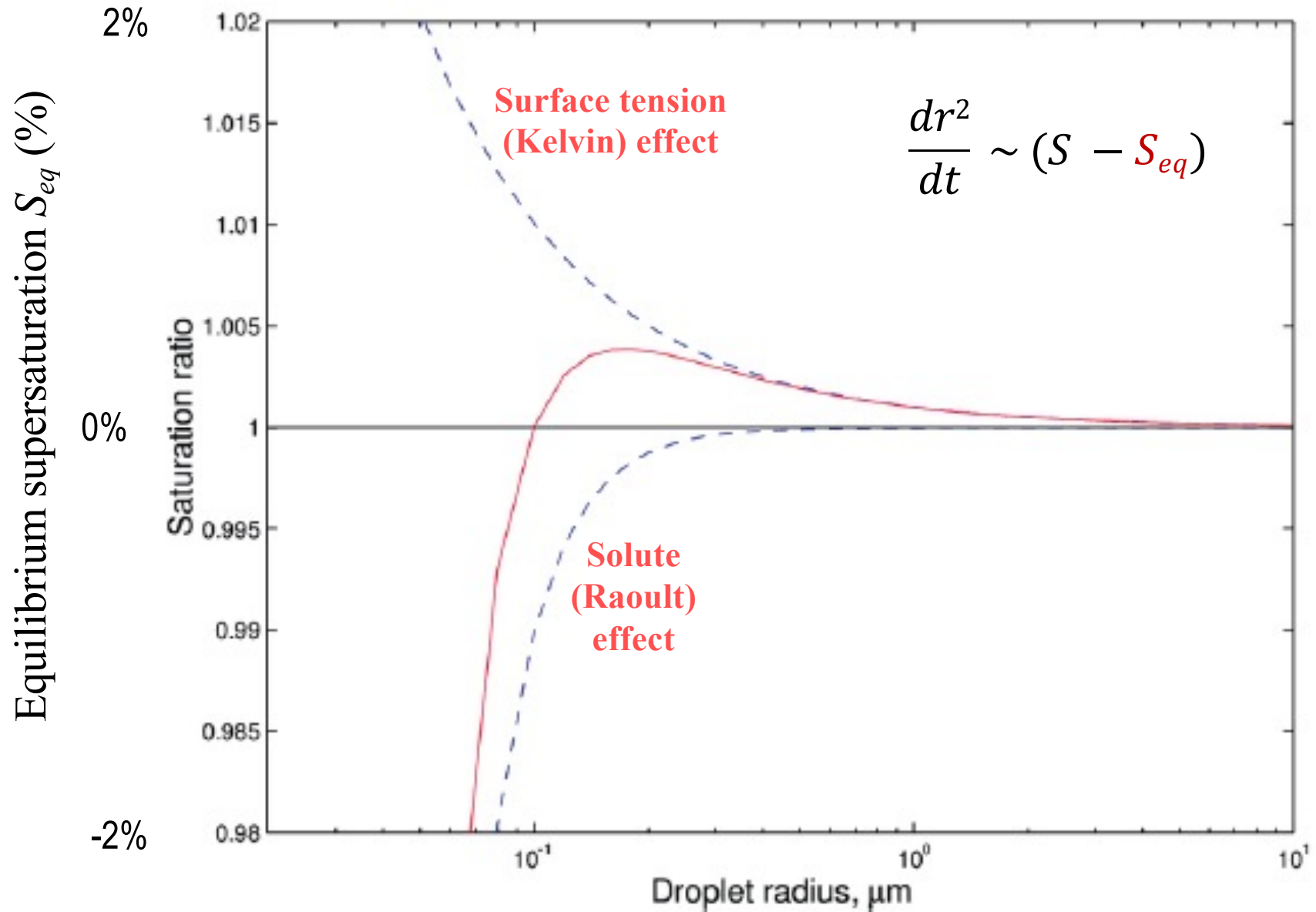
G – temperature and pressure dependent coefficient

$$S_{eq} = A/r - B/r^3$$

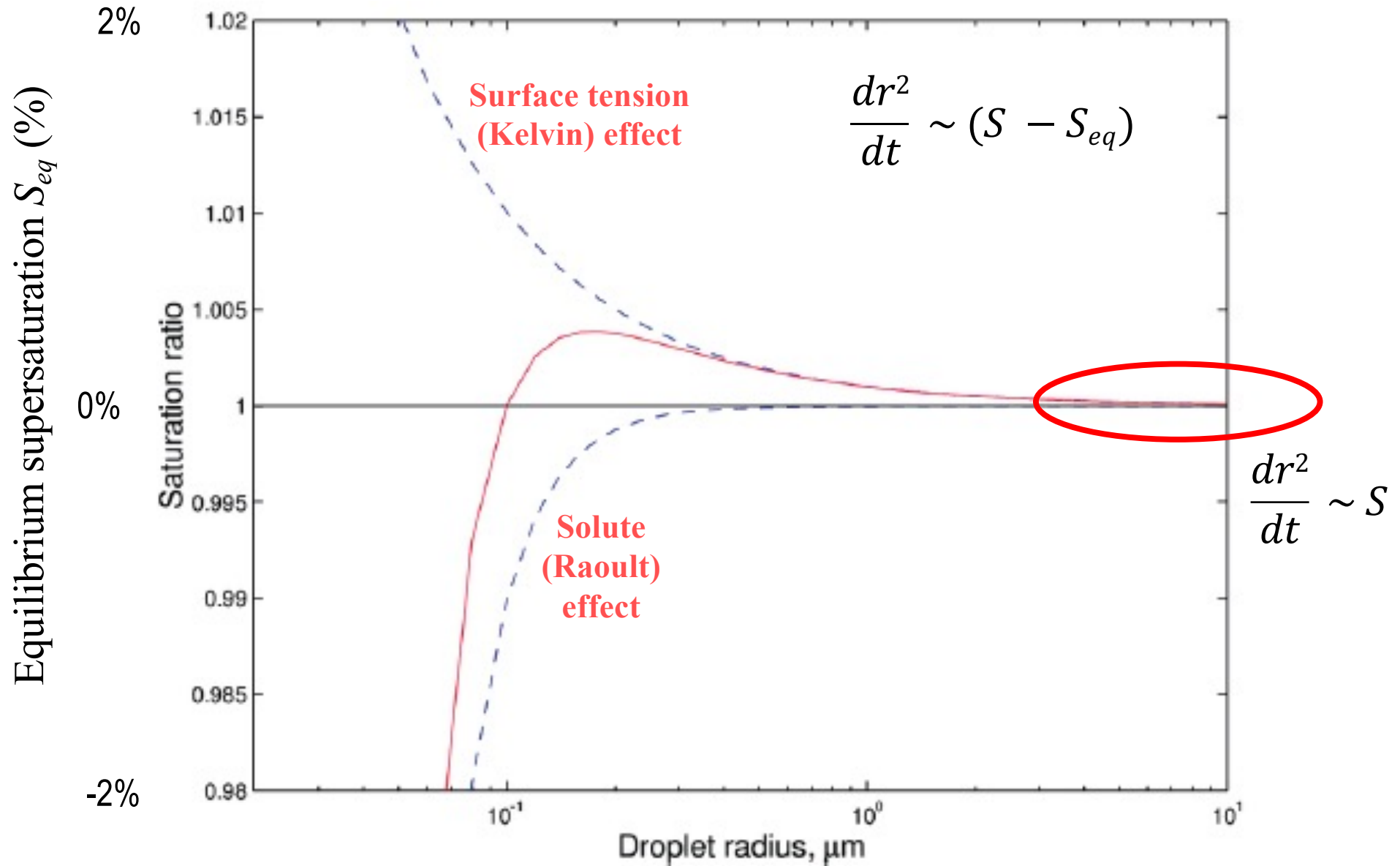
A represents Kelvin or curvature effect;

B depends on the solute properties and represents Raoult effect

Equilibrium supersaturation S_{eq} as a function of a droplet radius for a given CCN size

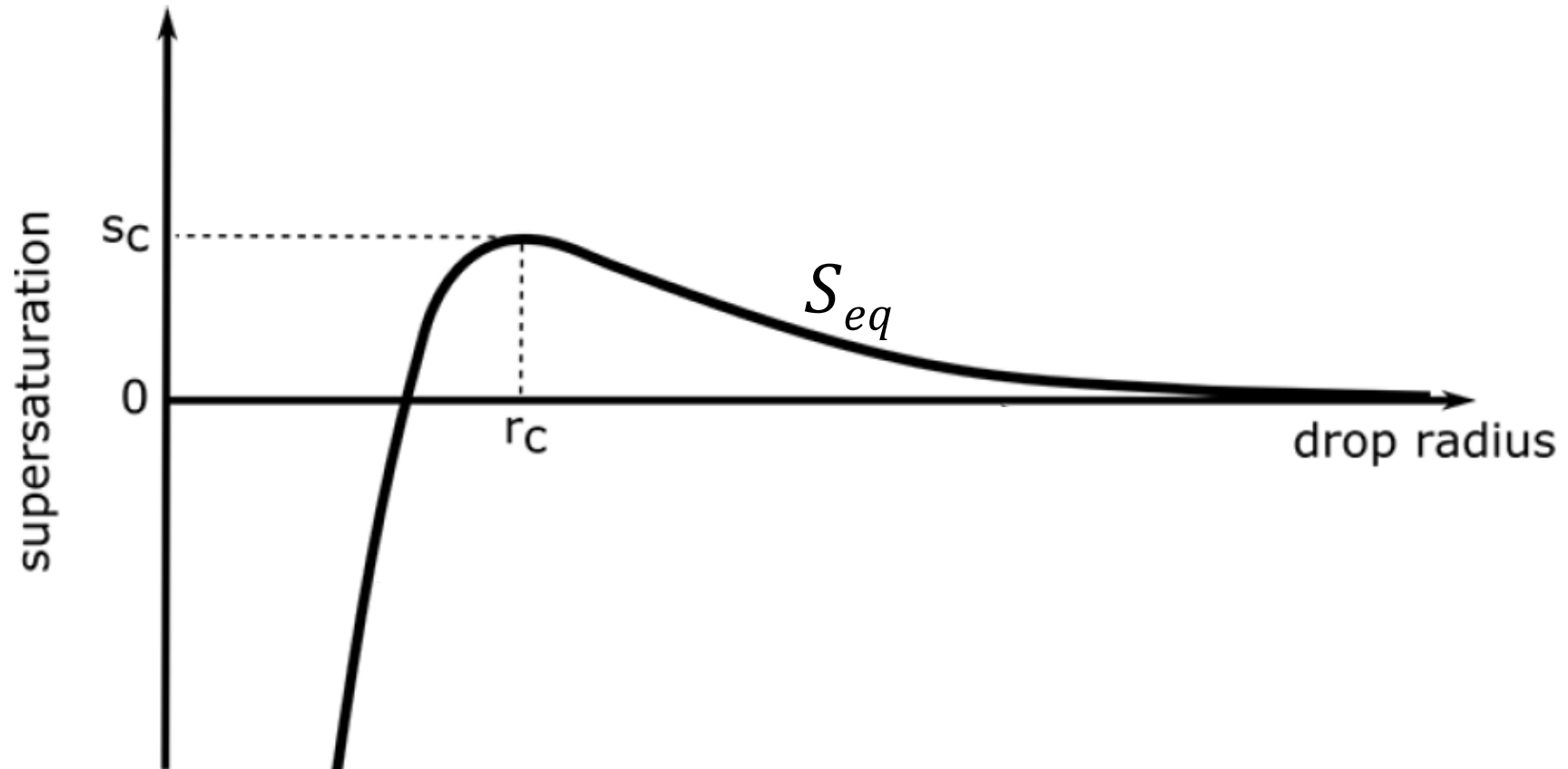


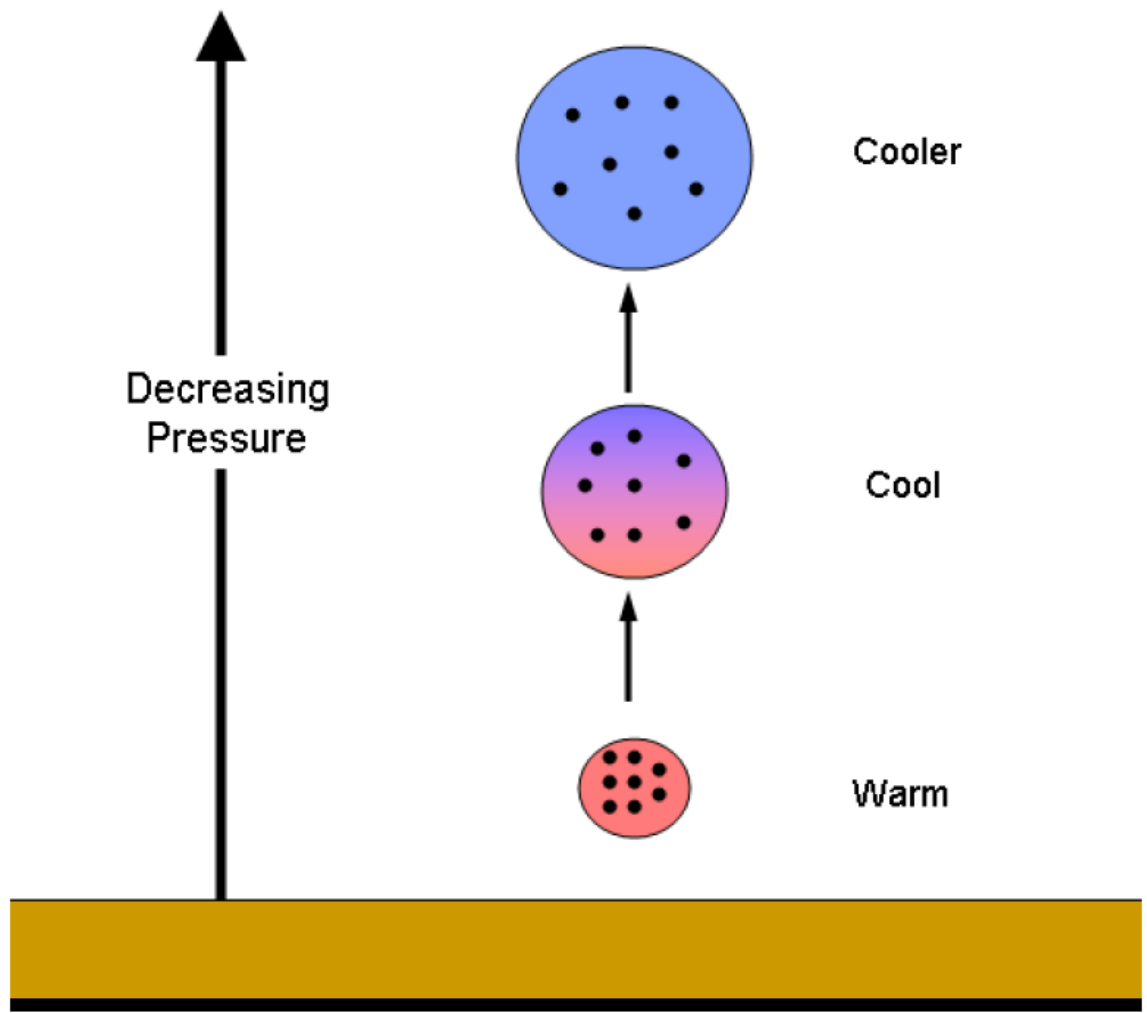
Equilibrium supersaturation S_{eq} as a function of a droplet radius



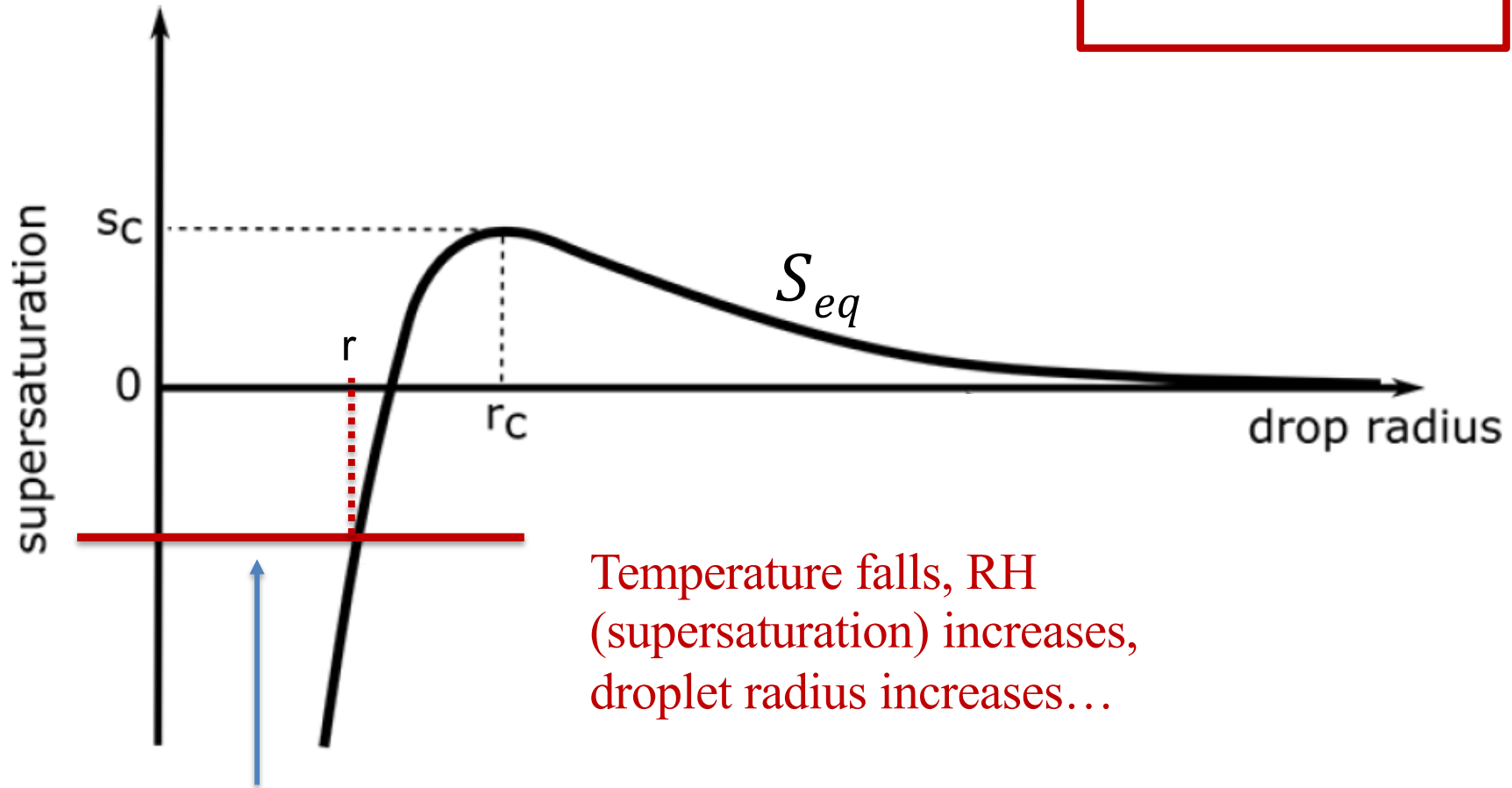
Droplet growth or evaporation depends on the difference between supersaturation S and equilibrium supersaturation S_{eq} for a given droplet radius:

$$\frac{dr^2}{dt} \sim (S - S_{eq})$$





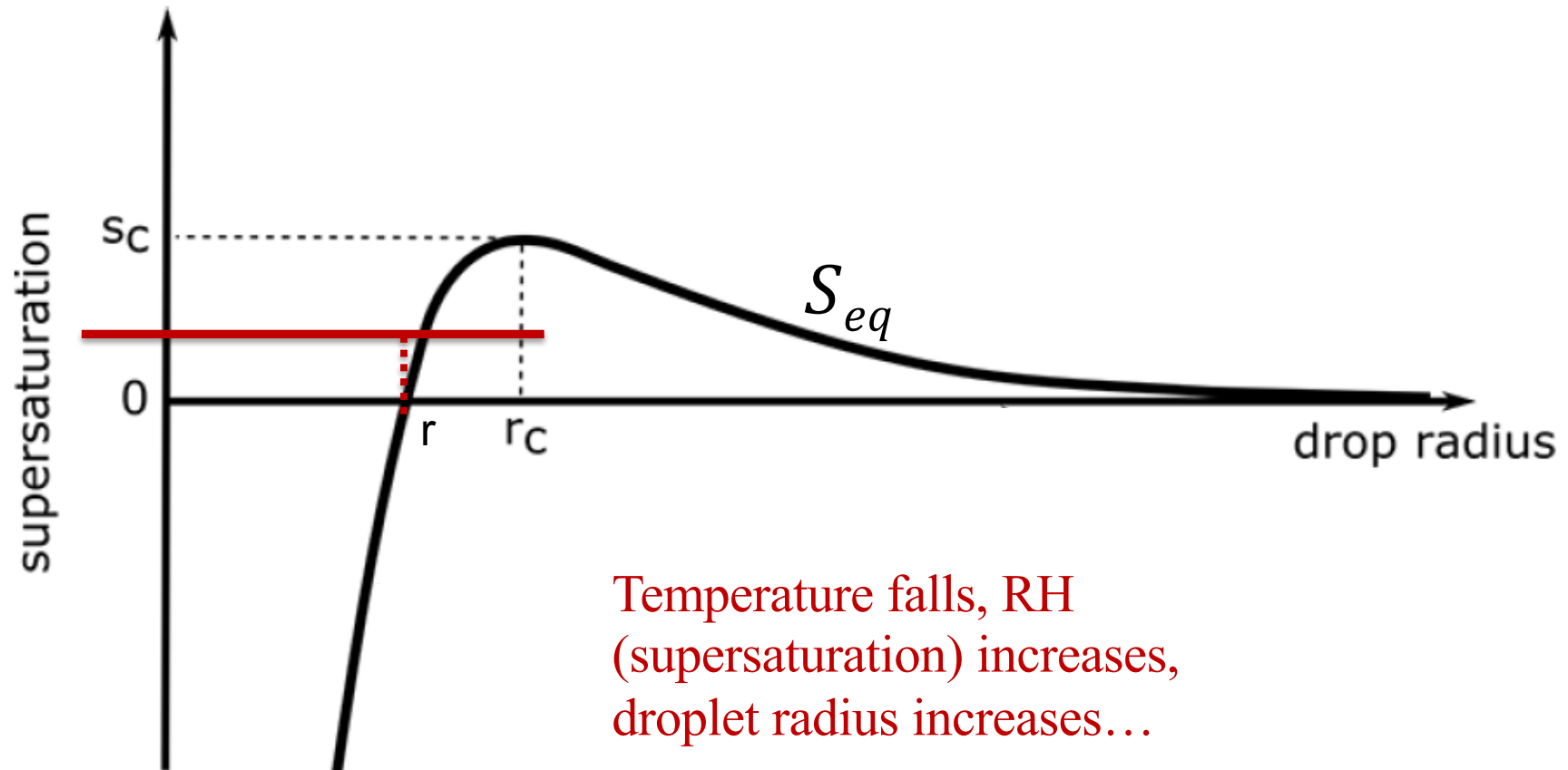
$$\frac{dr^2}{dt} \sim (S - S_{eq})$$



Temperature falls, RH (supersaturation) increases, droplet radius increases...

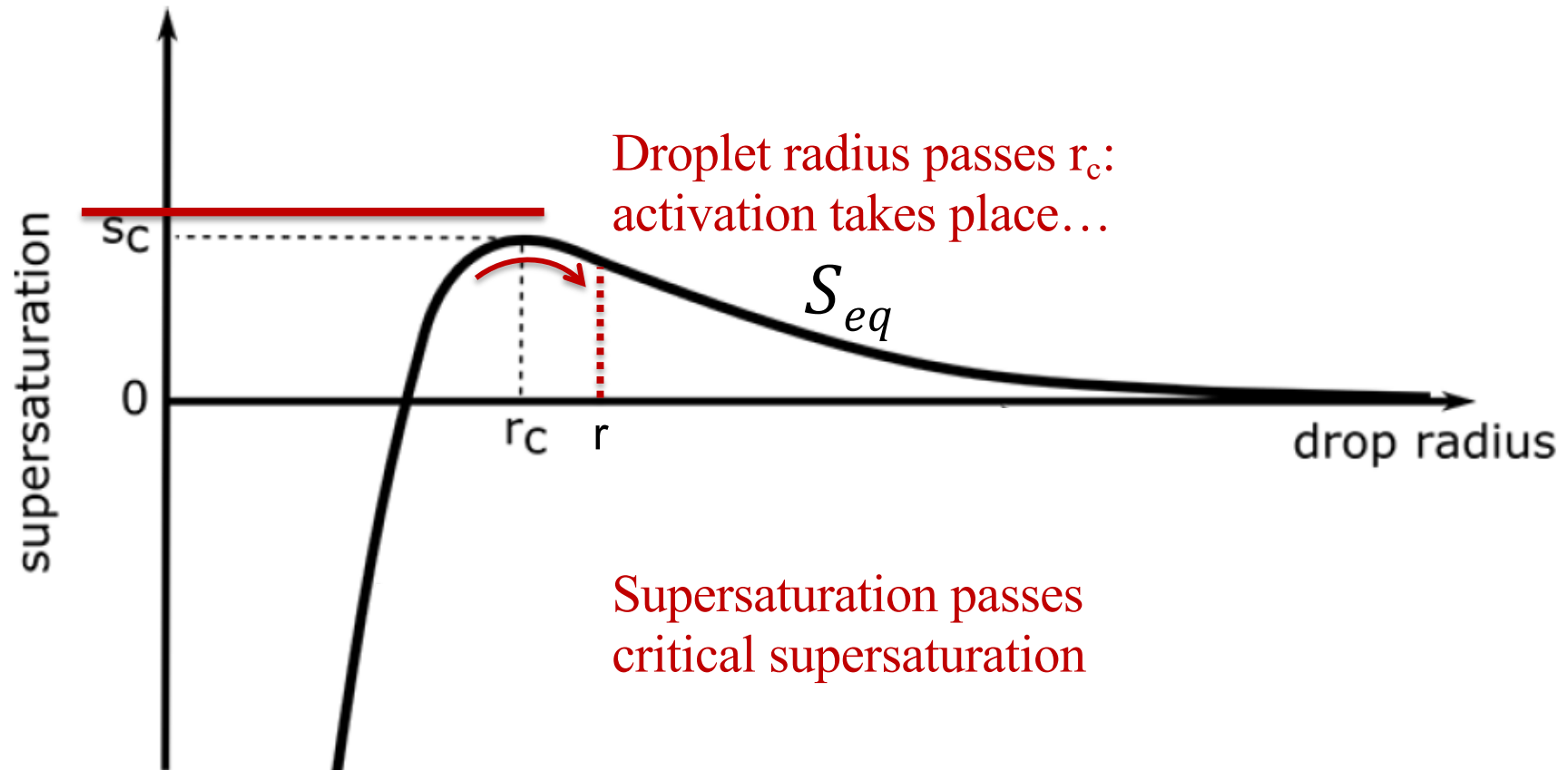
This line shows environmental supersaturation (or RH)

$$\frac{dr^2}{dt} \sim (S - S_{eq})$$

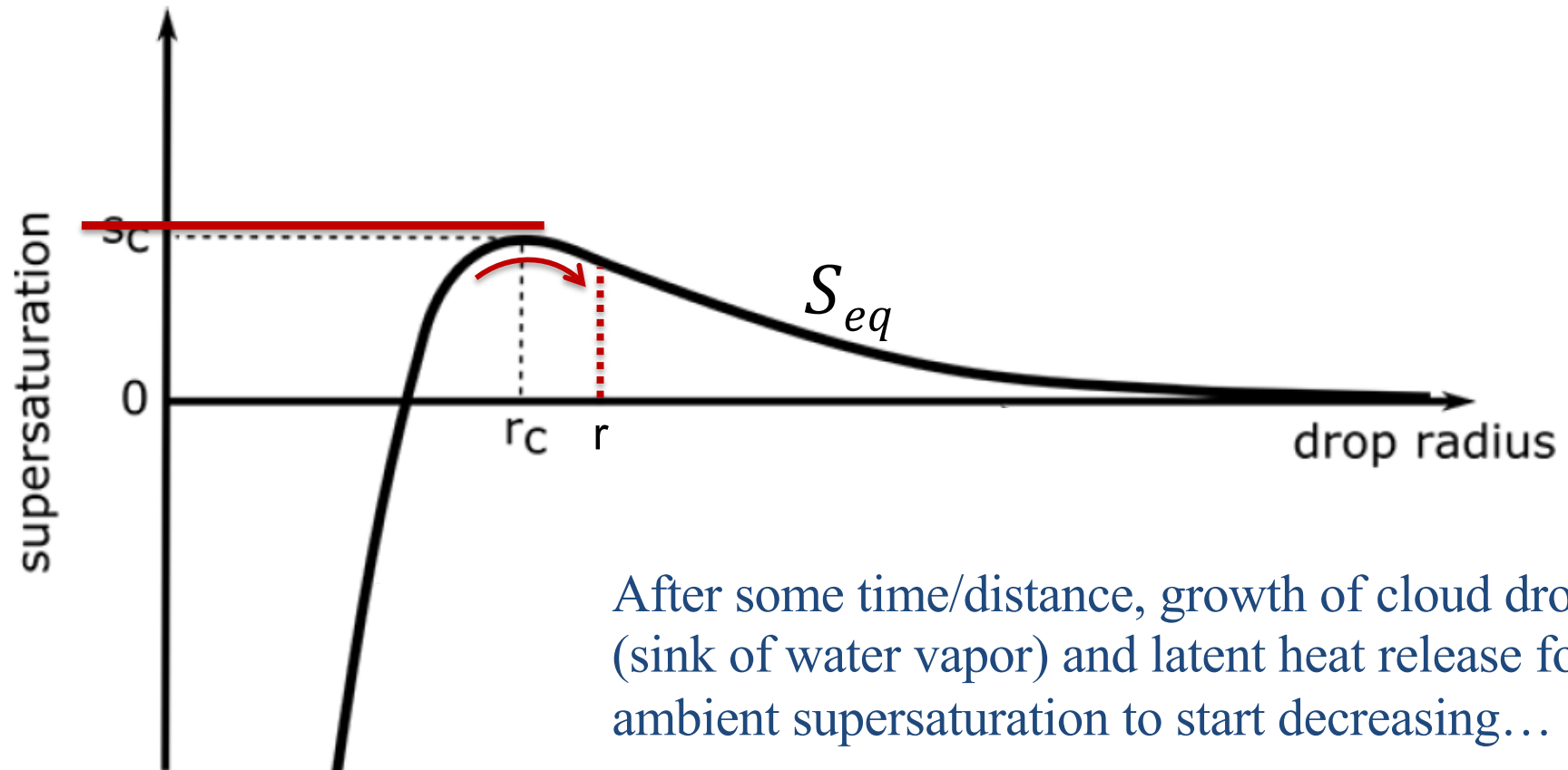


Temperature falls, RH
(supersaturation) increases,
droplet radius increases...

$$\frac{dr^2}{dt} \sim (S - S_{eq})$$

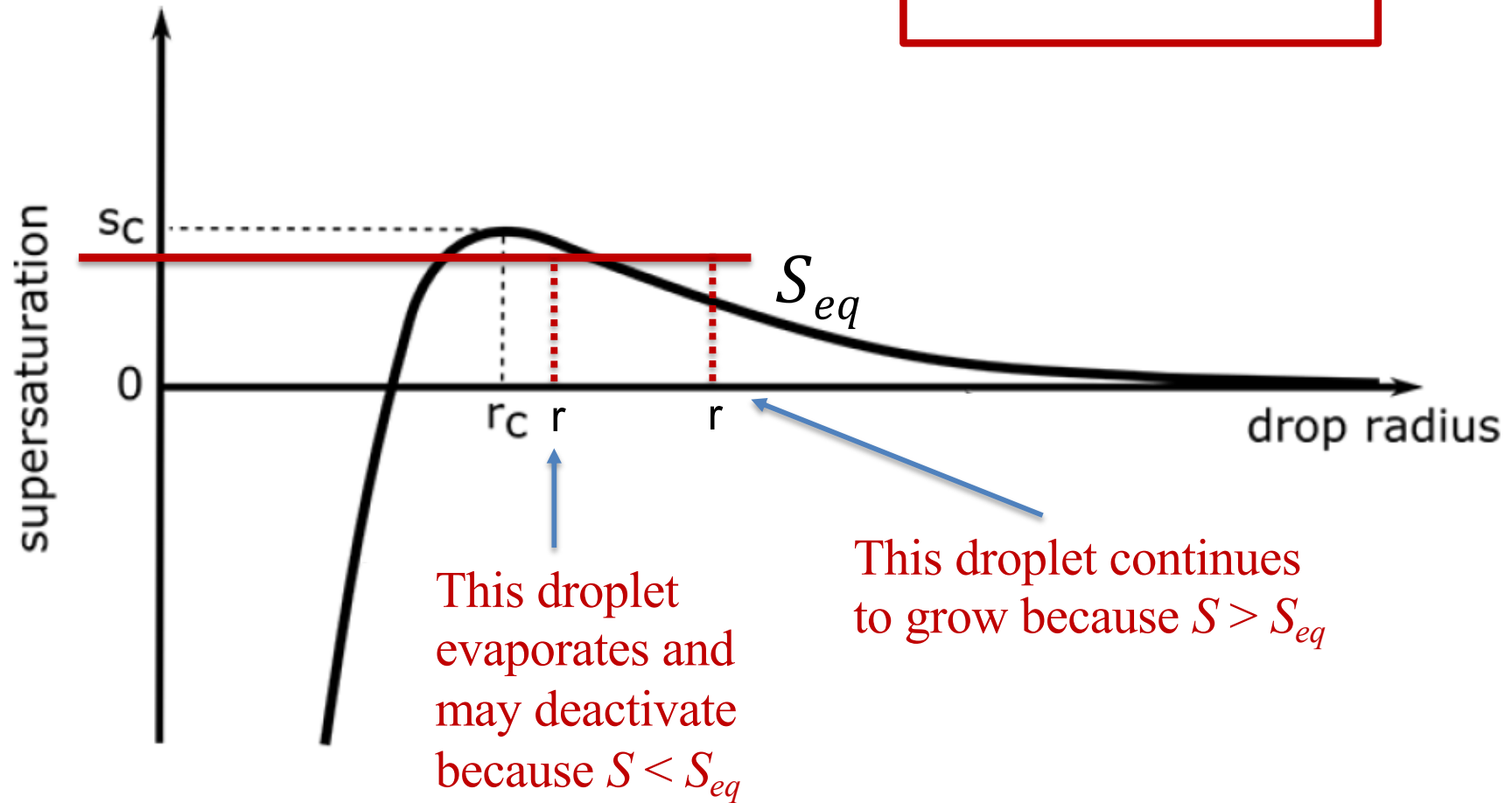


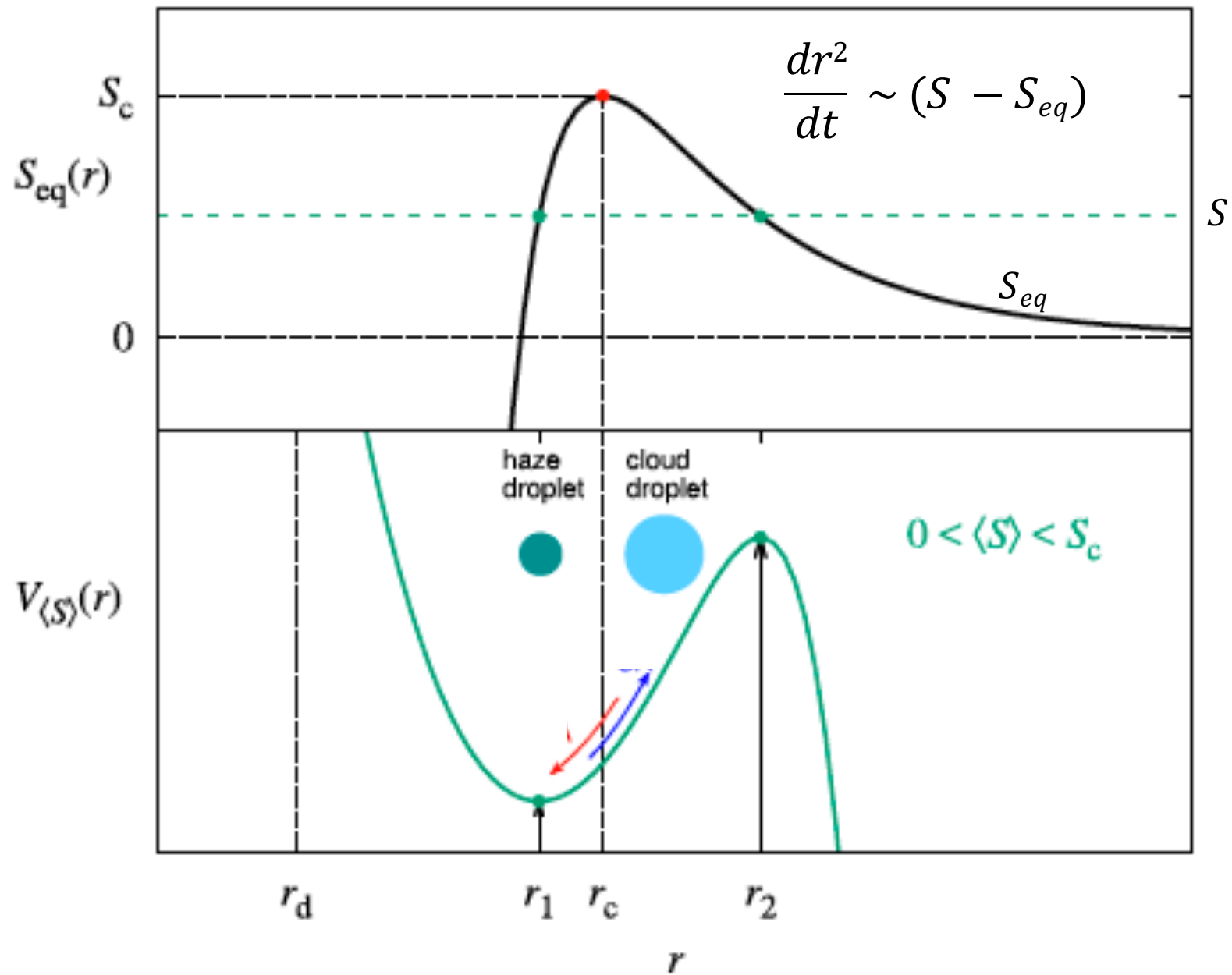
$$\frac{dr^2}{dt} \sim (S - S_{eq})$$



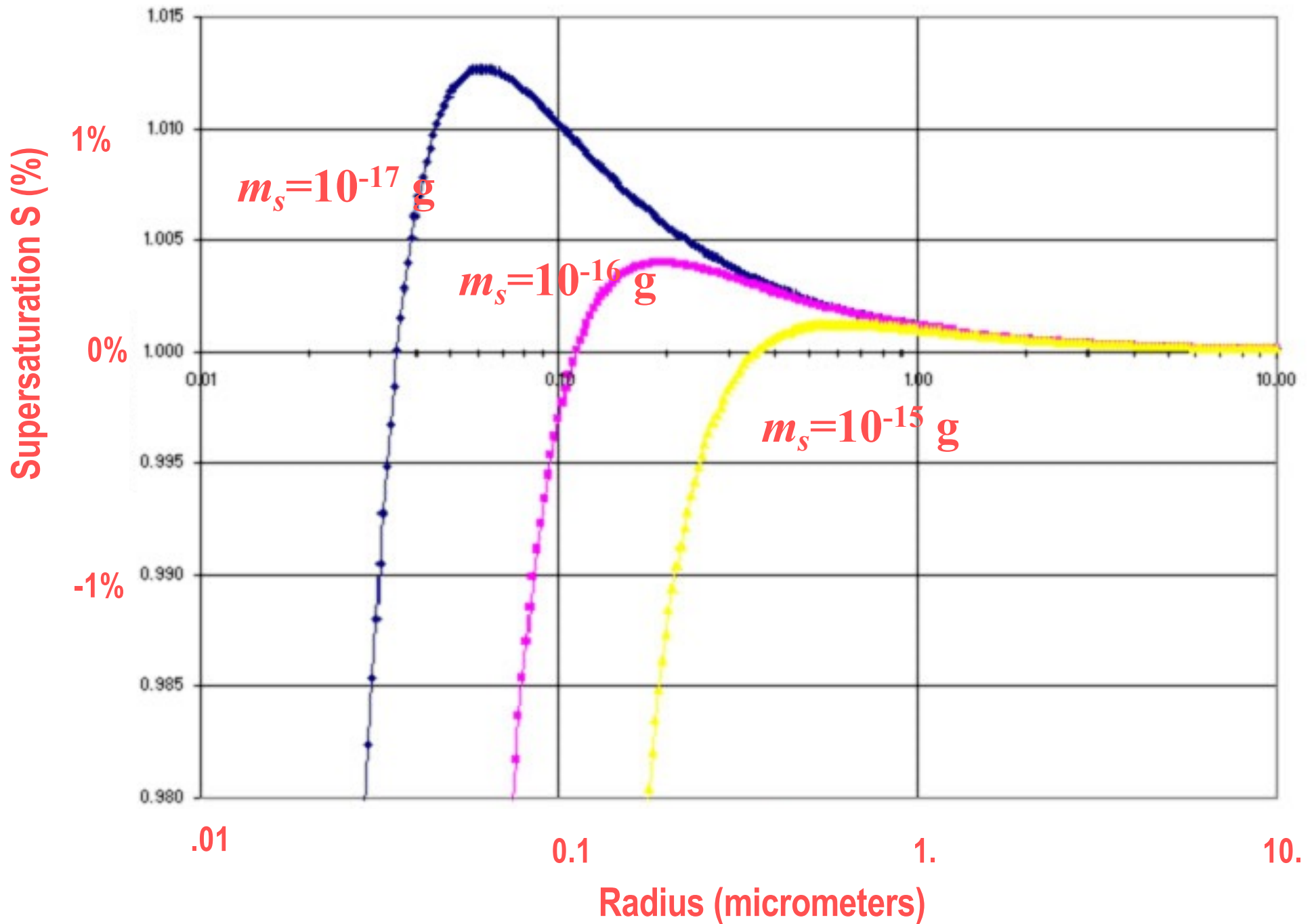
After some time/distance, growth of cloud droplets (sink of water vapor) and latent heat release forces ambient supersaturation to start decreasing...

$$\frac{dr^2}{dt} \sim (S - S_{eq})$$



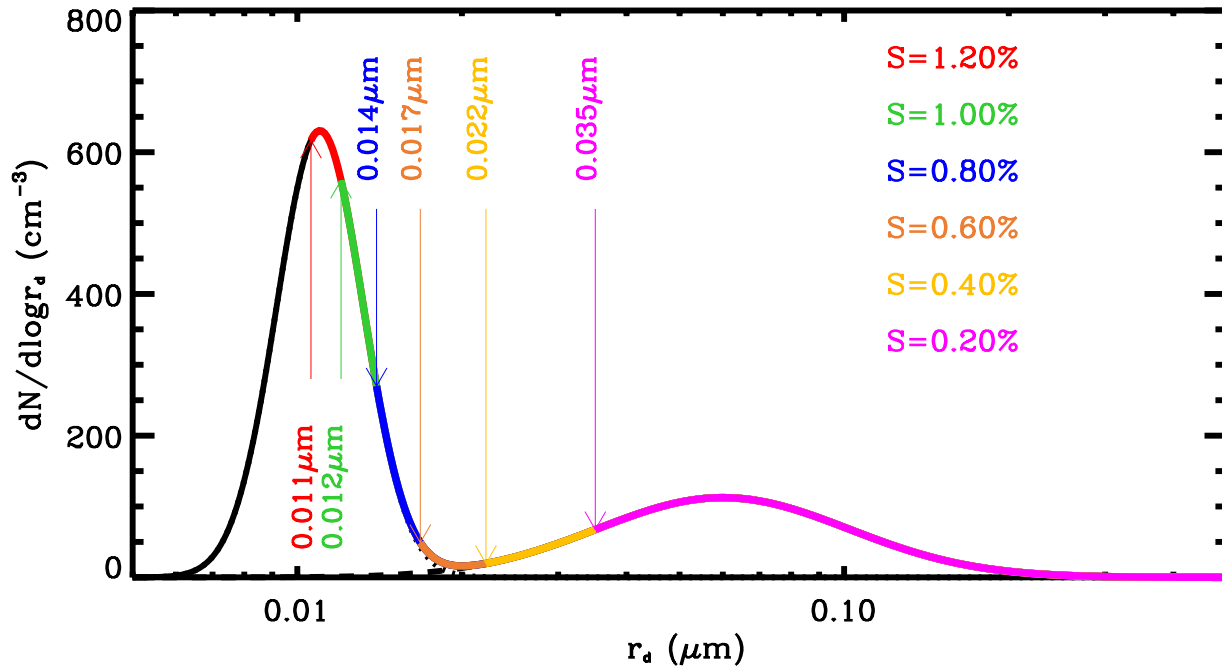


Köhler curves for different CCN dry mass m_s (or dry radius):



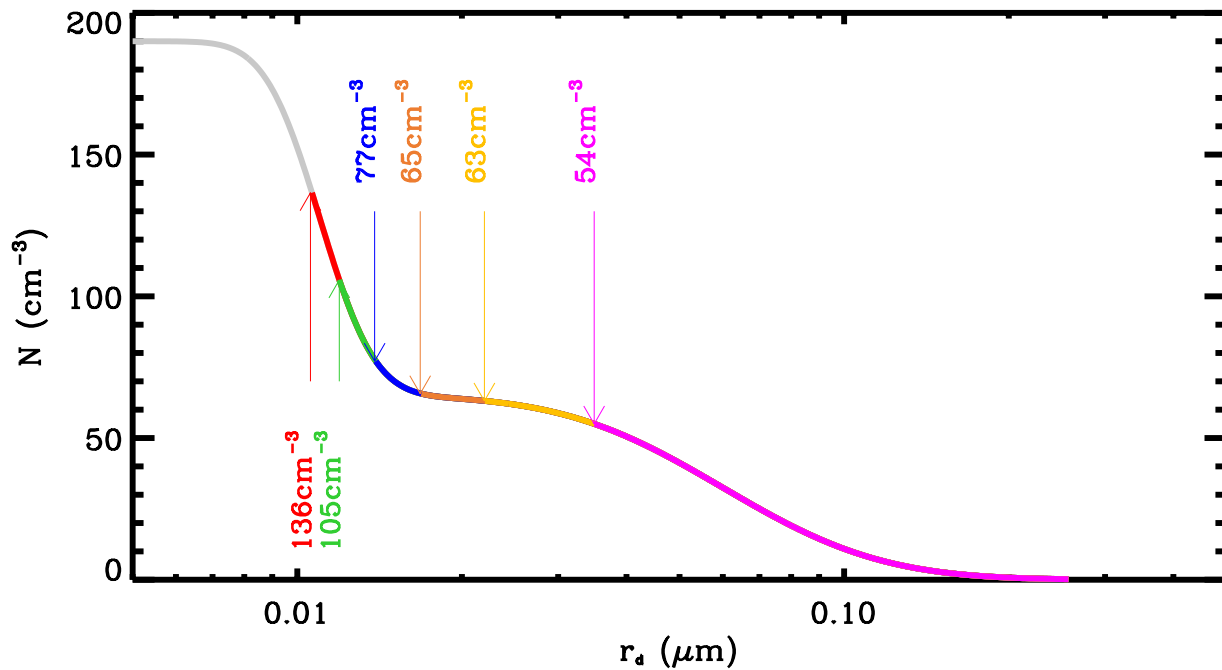
The role of CCN distribution: contrast **pristine** and **polluted** environments

marine aerosol, $\kappa_{\text{mean}}=1.28$

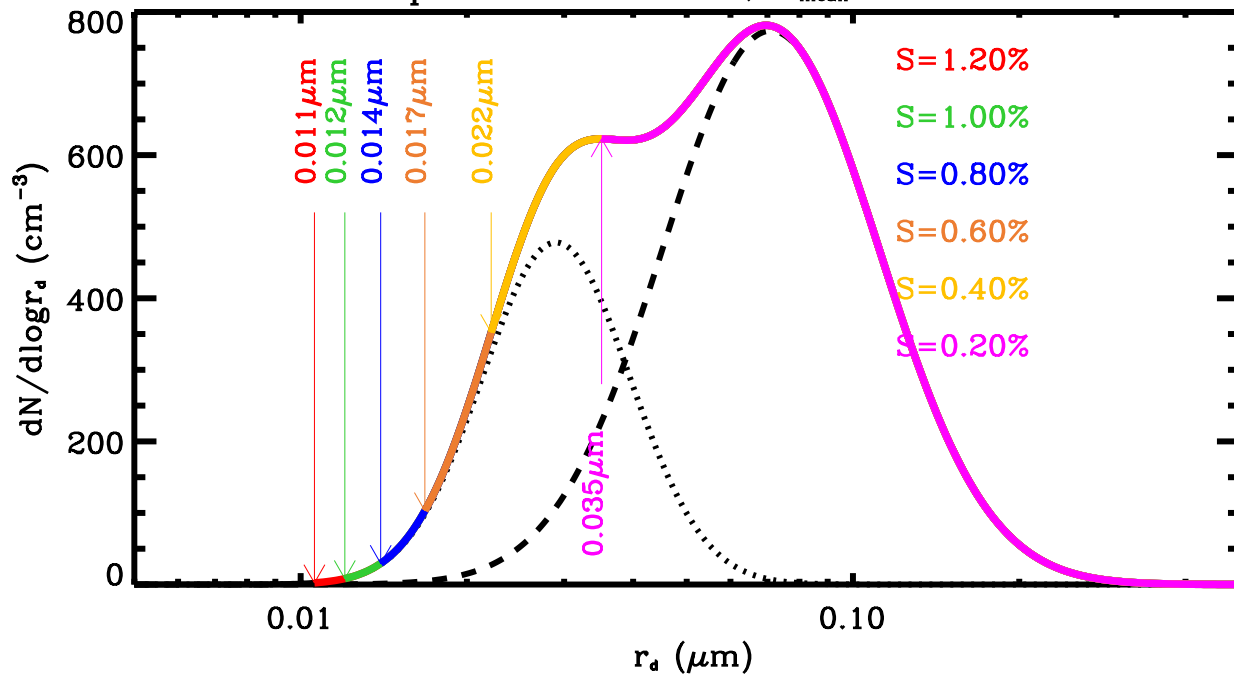


Pristine (PRI) or maritime aerosol from DYCOMS observations (Pacific off California coast, Grabowski et al. *Atmos. Res.* 2011).

marine aerosol

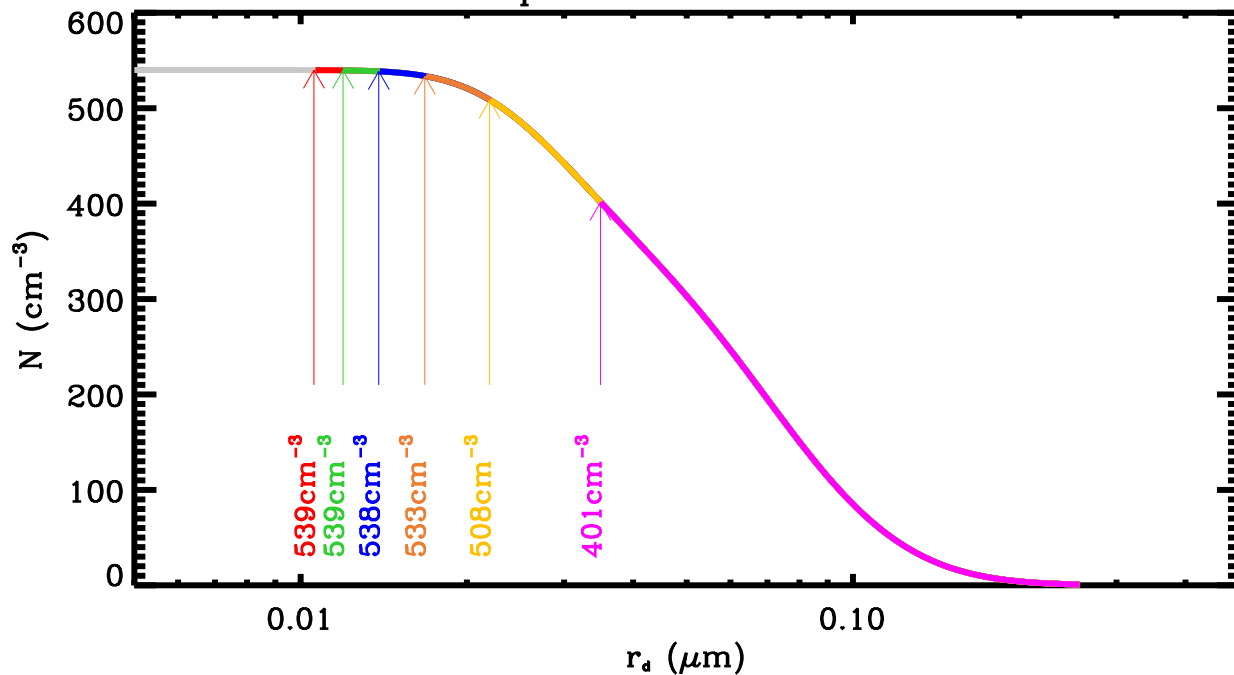


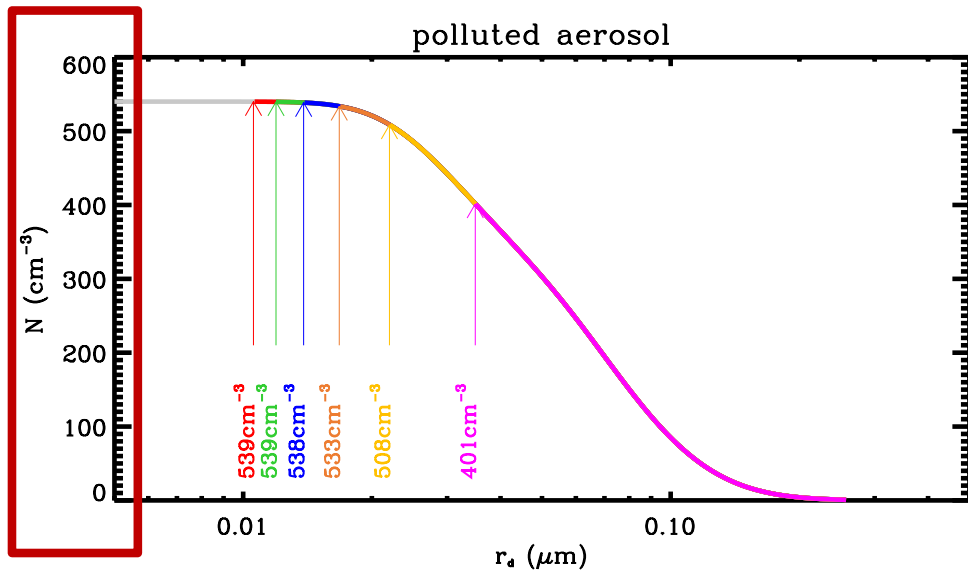
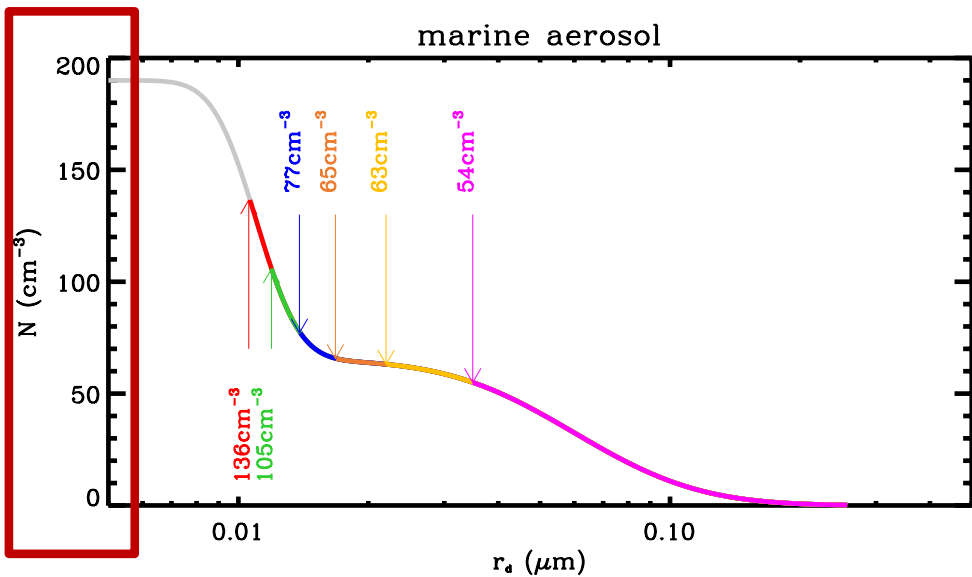
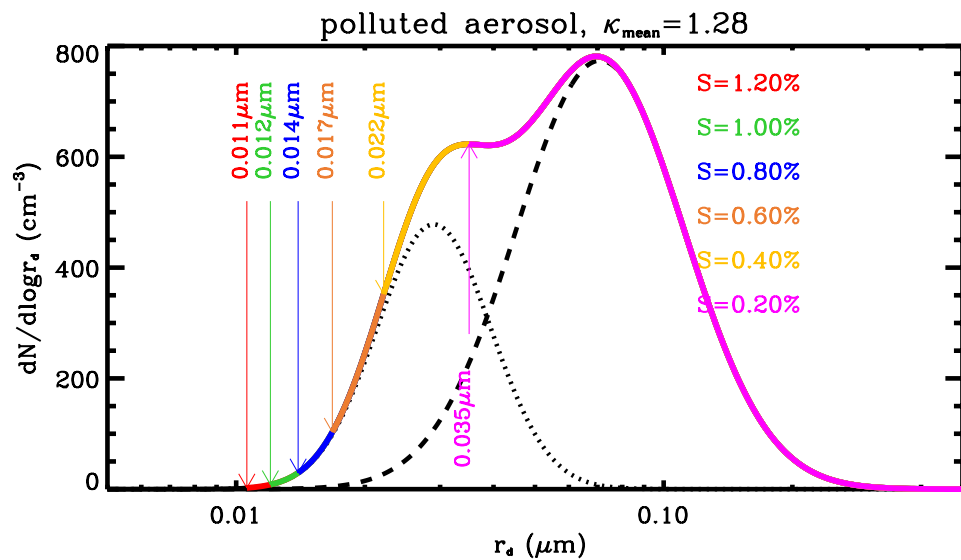
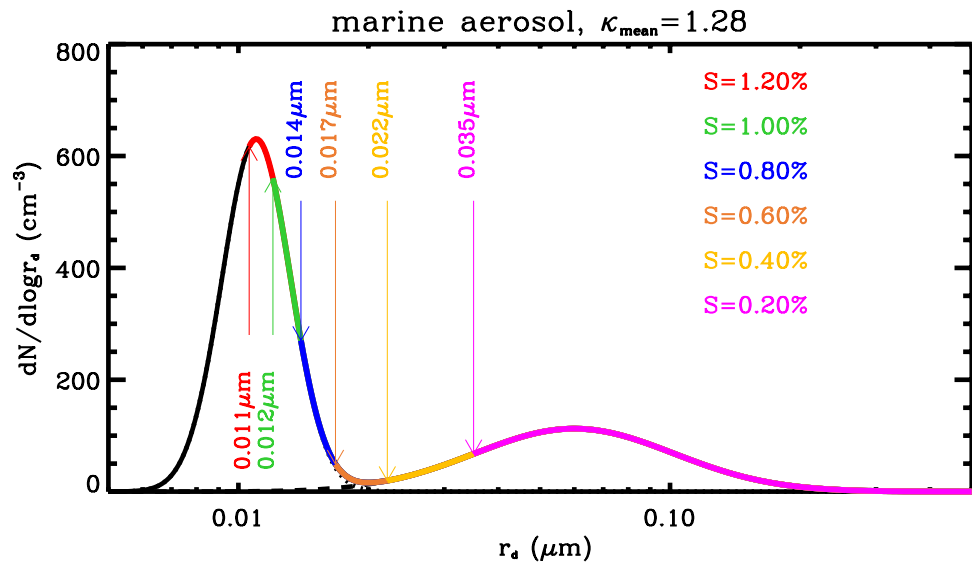
polluted aerosol, $\kappa_{\text{mean}}=1.28$



Polluted (POL) aerosol from VOCALS observations (S. Pacific off Chile coast, Grabowski et al. *Atmos. Res.* 2011).

polluted aerosol





adiabatic parcel crossing the cloud base:

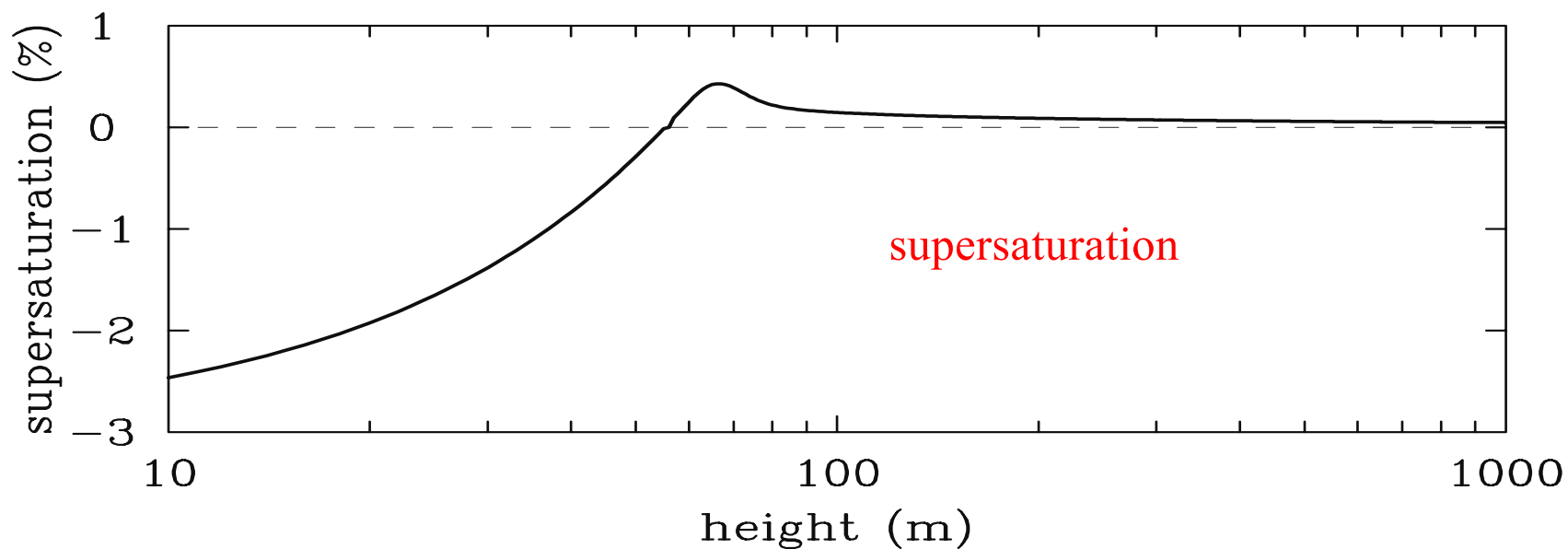
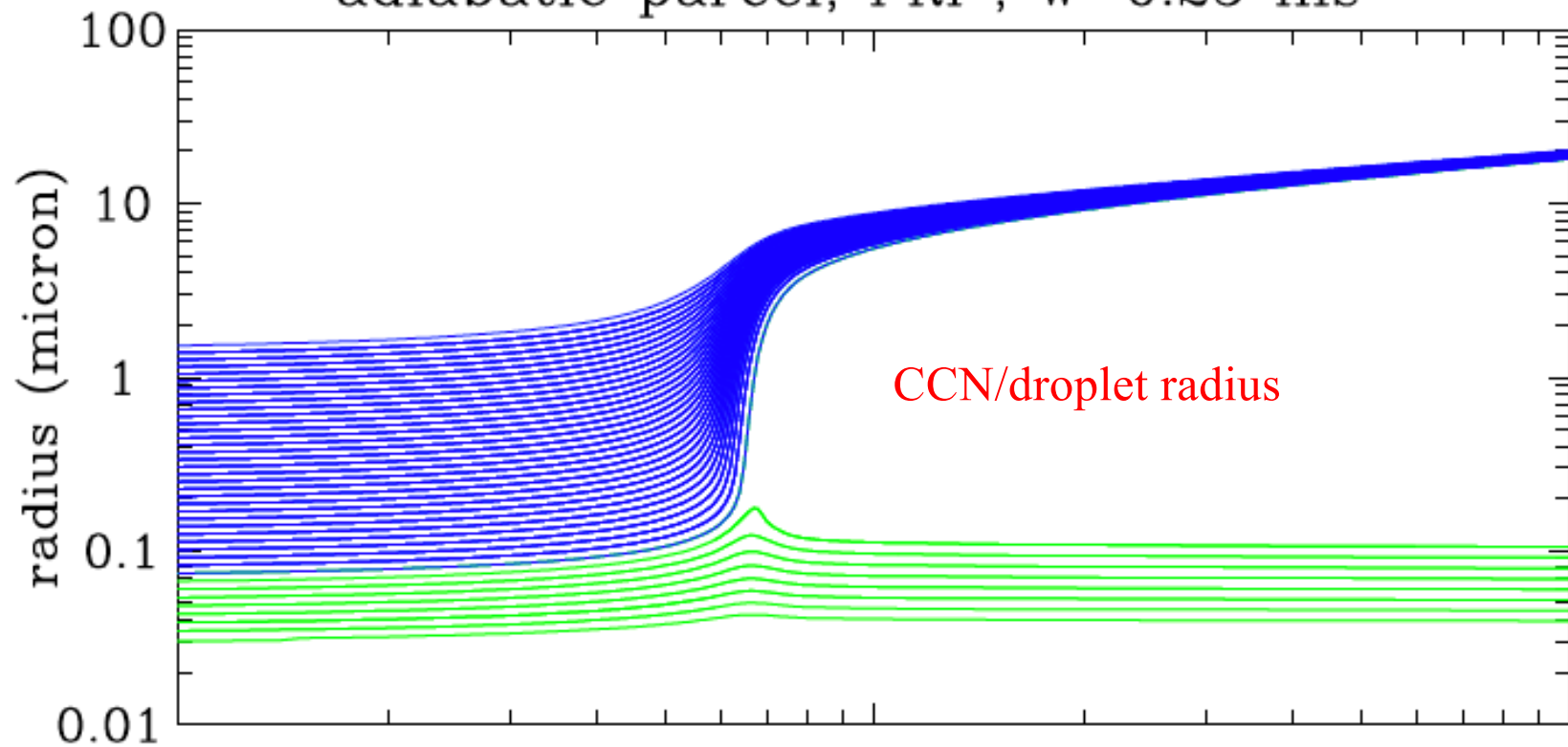
Starting from $RH = 97\%$, $p = 900$ hPa, $T = 282$ K

CCN (superdroplets) from a prescribed distribution
(assumed initially at equilibrium with ambient RH)

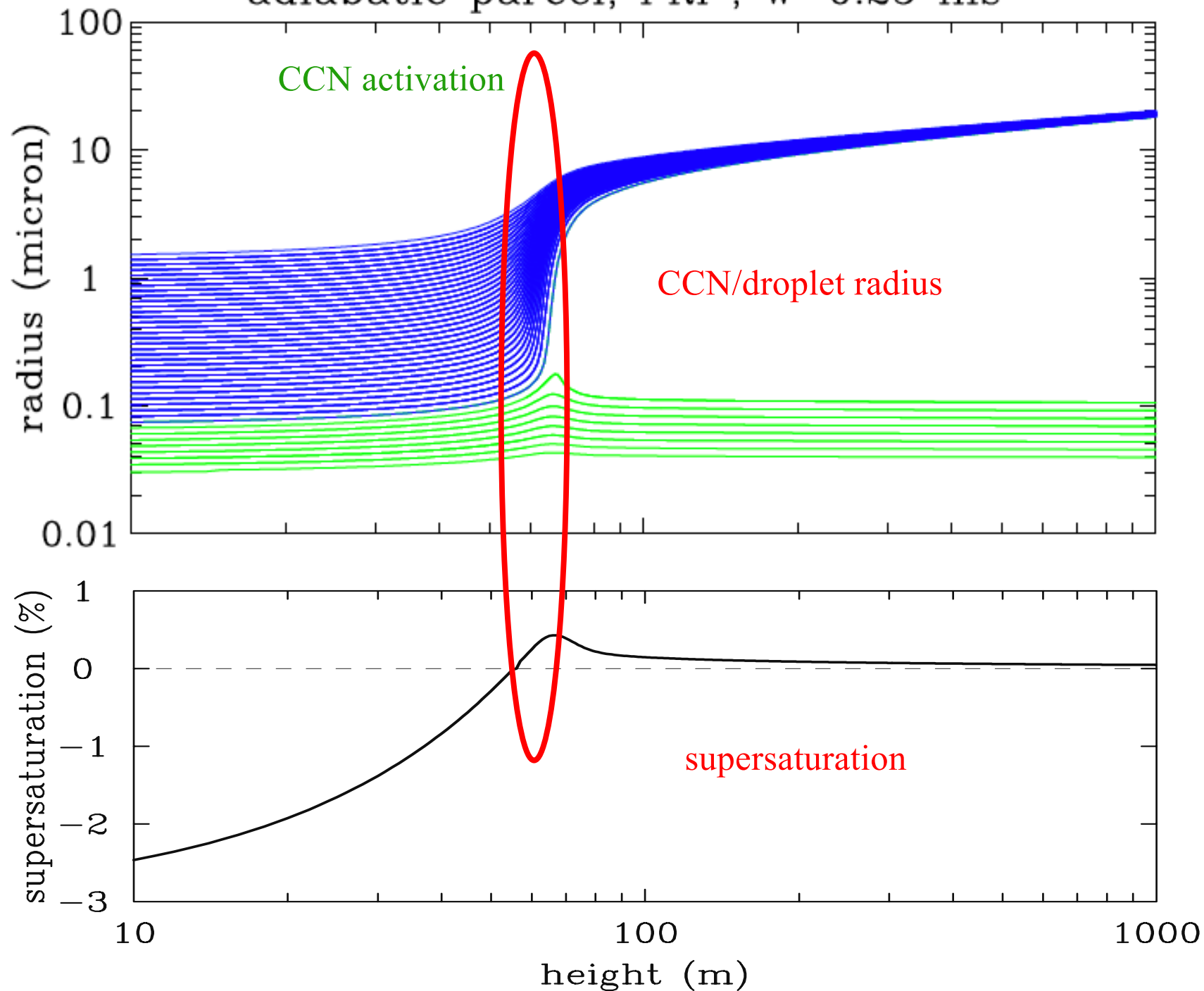
prescribed ascent of 0.25 , 1 , or 4 m s^{-1}

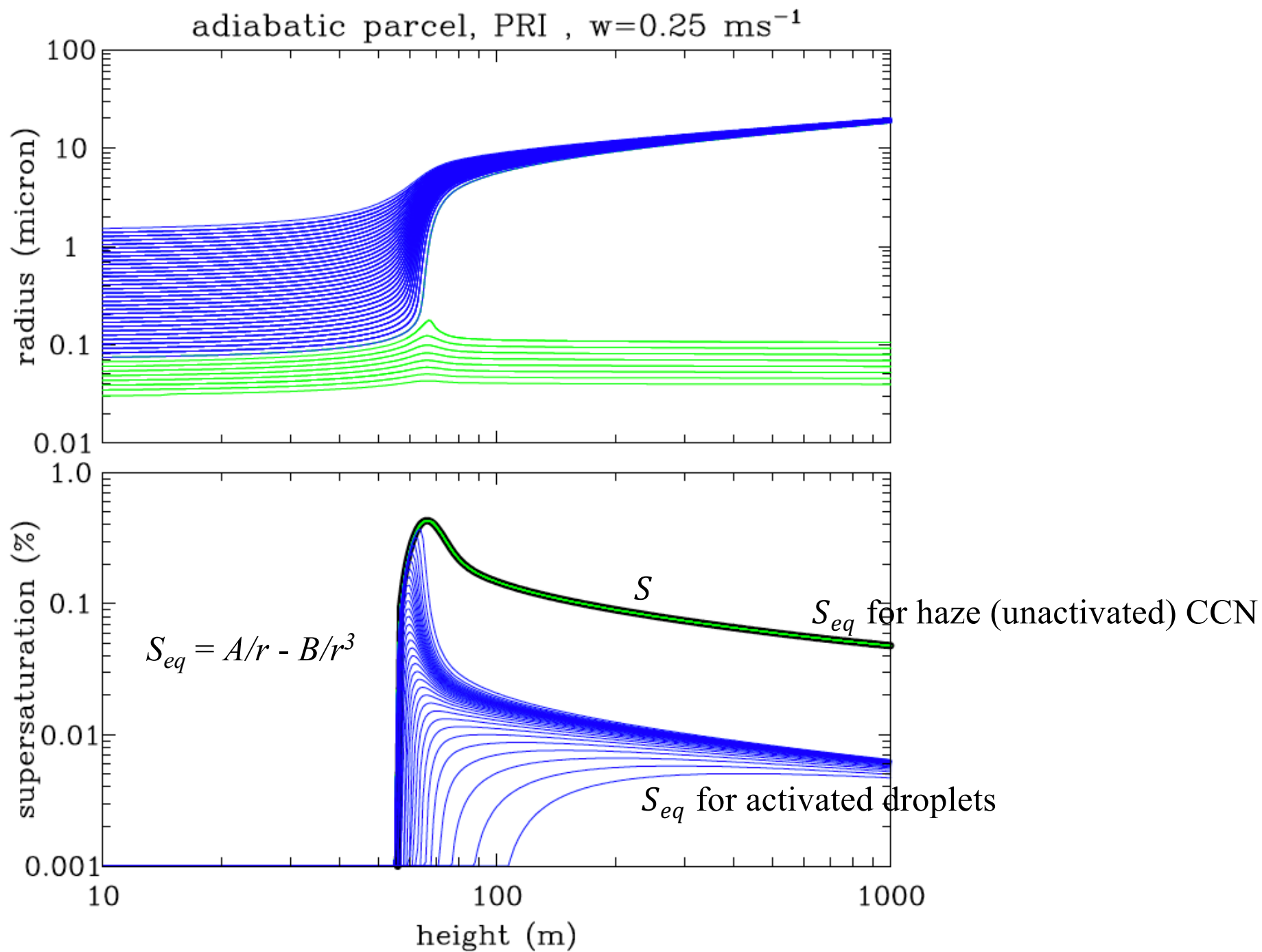
Parcel reaches saturation ($RH=100\%$) after about 50 m ascent;
we will look at the evolution of the spectral width up to 1 km.

adiabatic parcel, PRI , $w=0.25 \text{ ms}^{-1}$

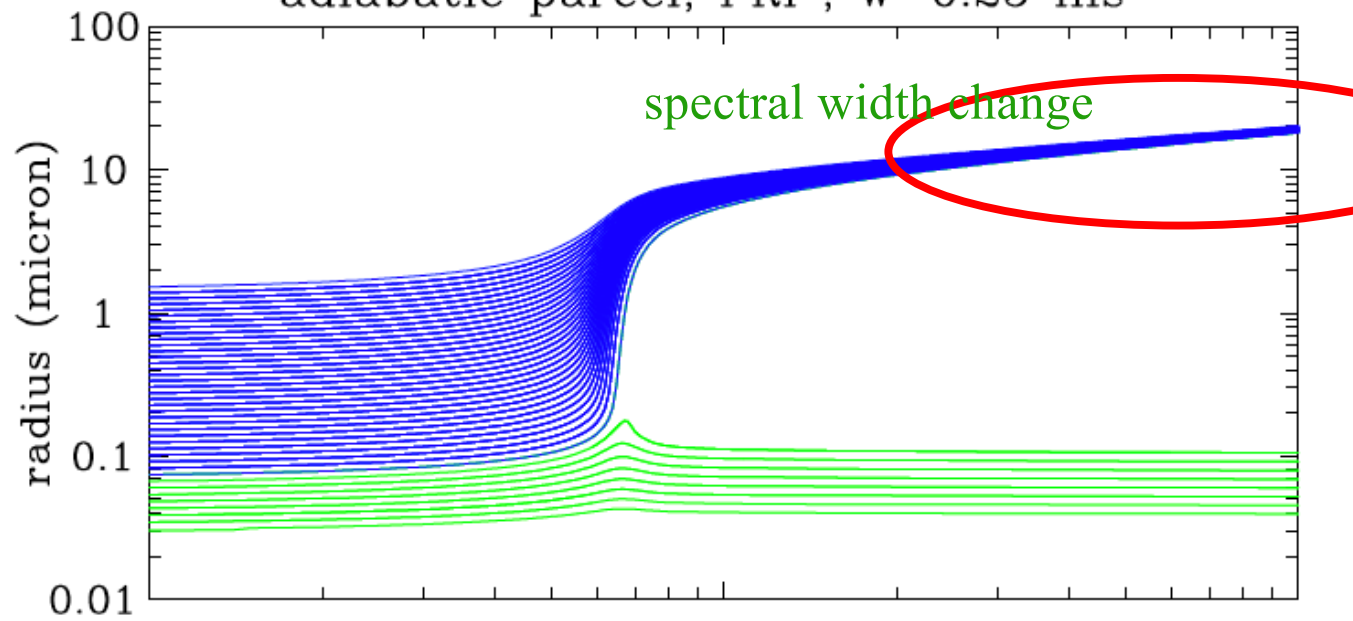


adiabatic parcel, PRI , $w=0.25 \text{ ms}^{-1}$





adiabatic parcel, PRI , $w=0.25 \text{ ms}^{-1}$

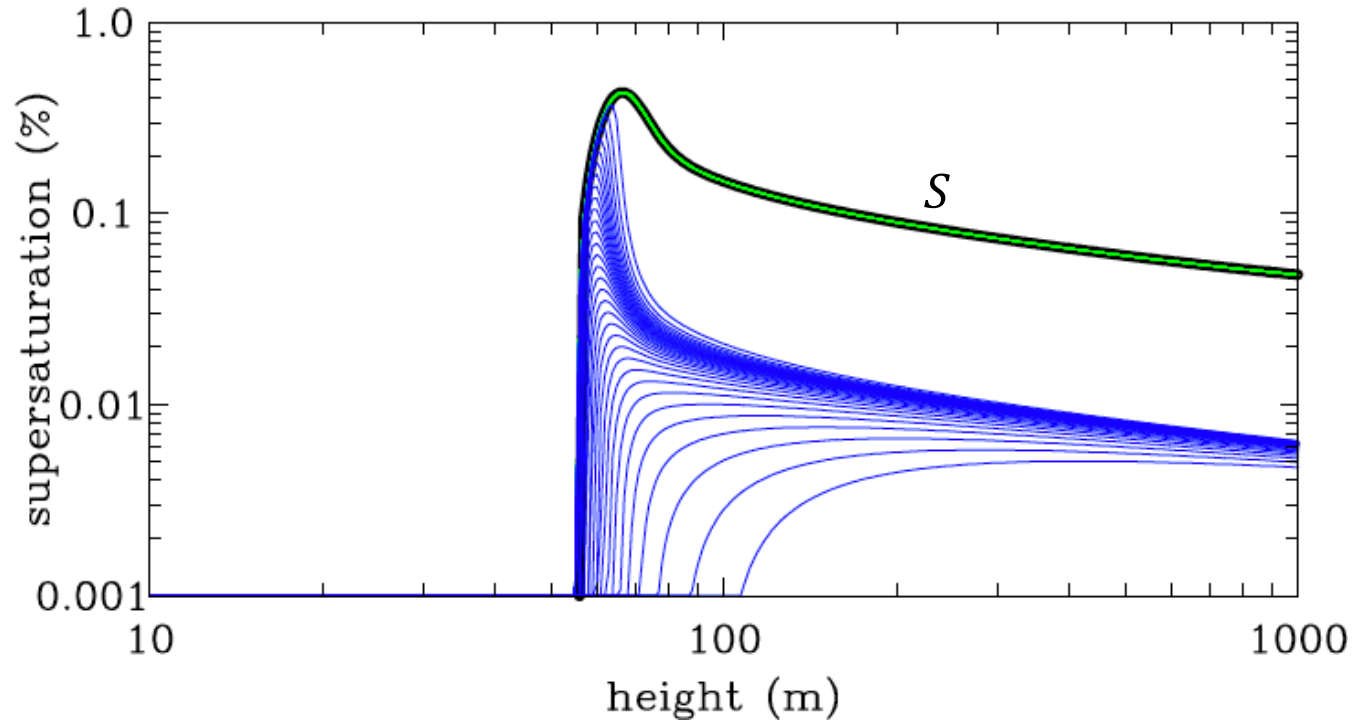


spectral width change

$$\frac{dr}{dt} \sim \frac{S}{r}$$

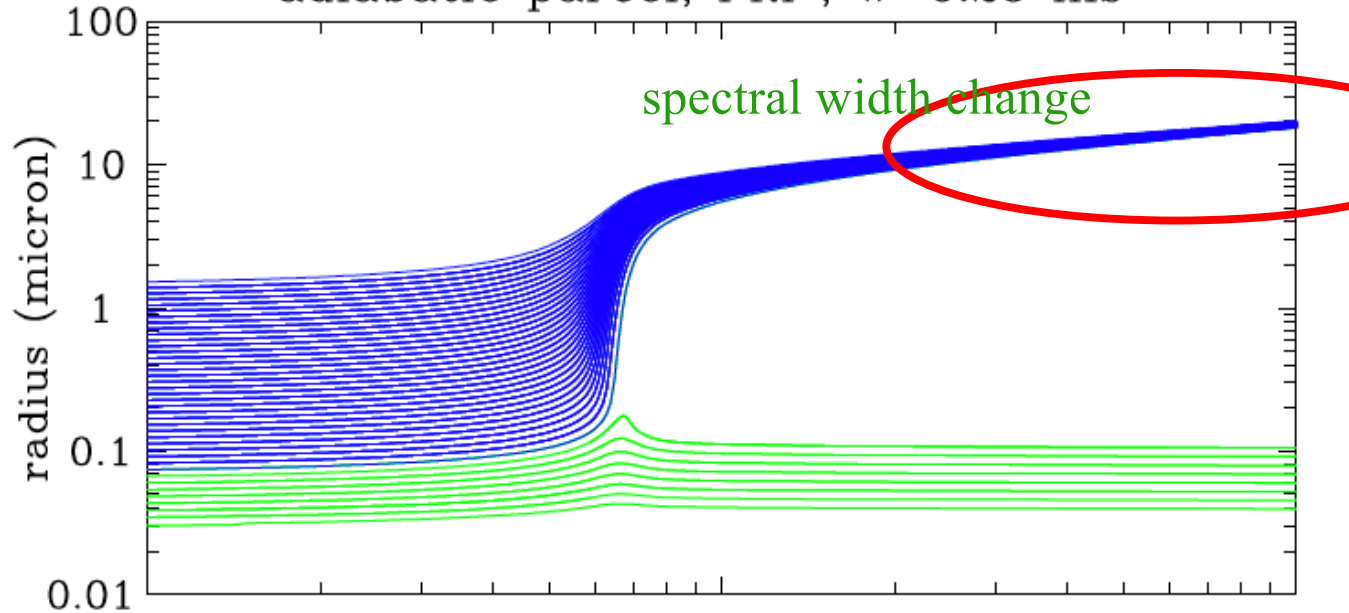
$$\frac{dr^2}{dt} \sim S$$

spectral width in r^2 space does not change when $S \gg S_{eq}$



S

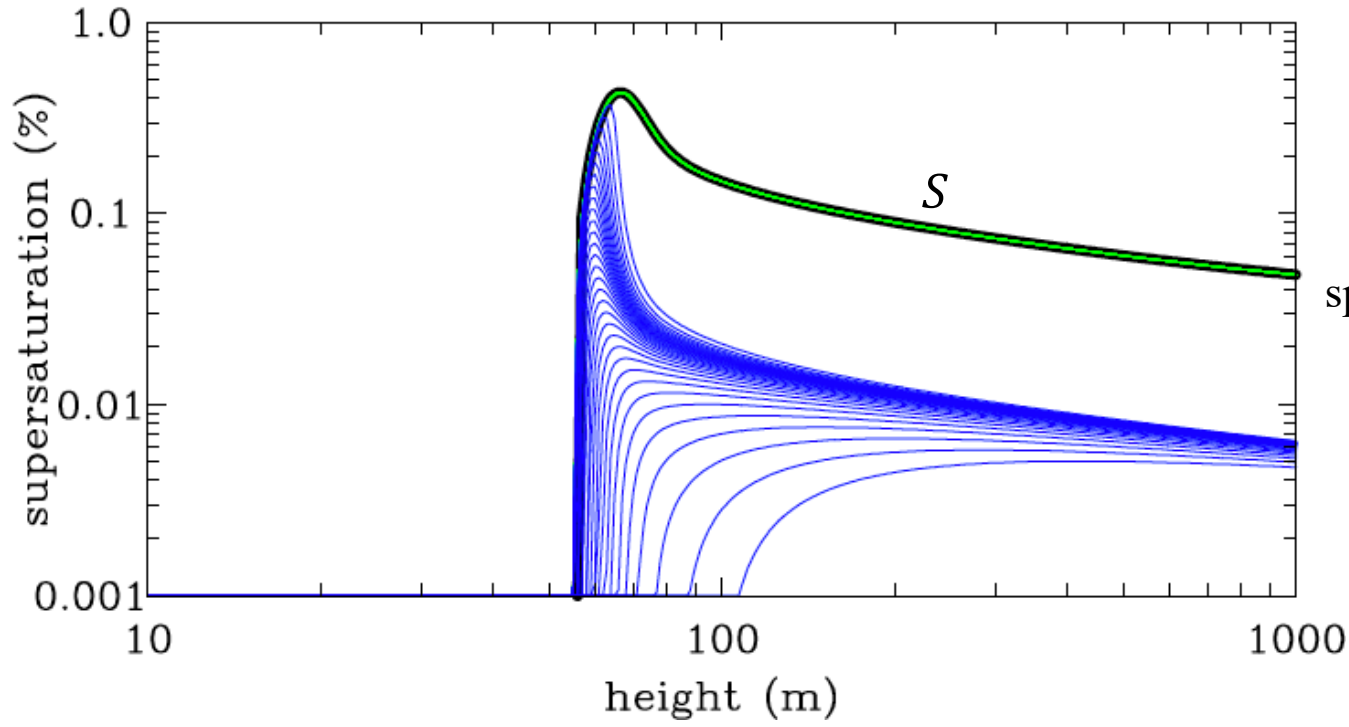
adiabatic parcel, PRI , $w=0.25 \text{ ms}^{-1}$



spectral width change

$$\frac{dr^2}{dt} \sim S$$

spectral width in r^2 space
does not change when $S \gg S_{eq}$

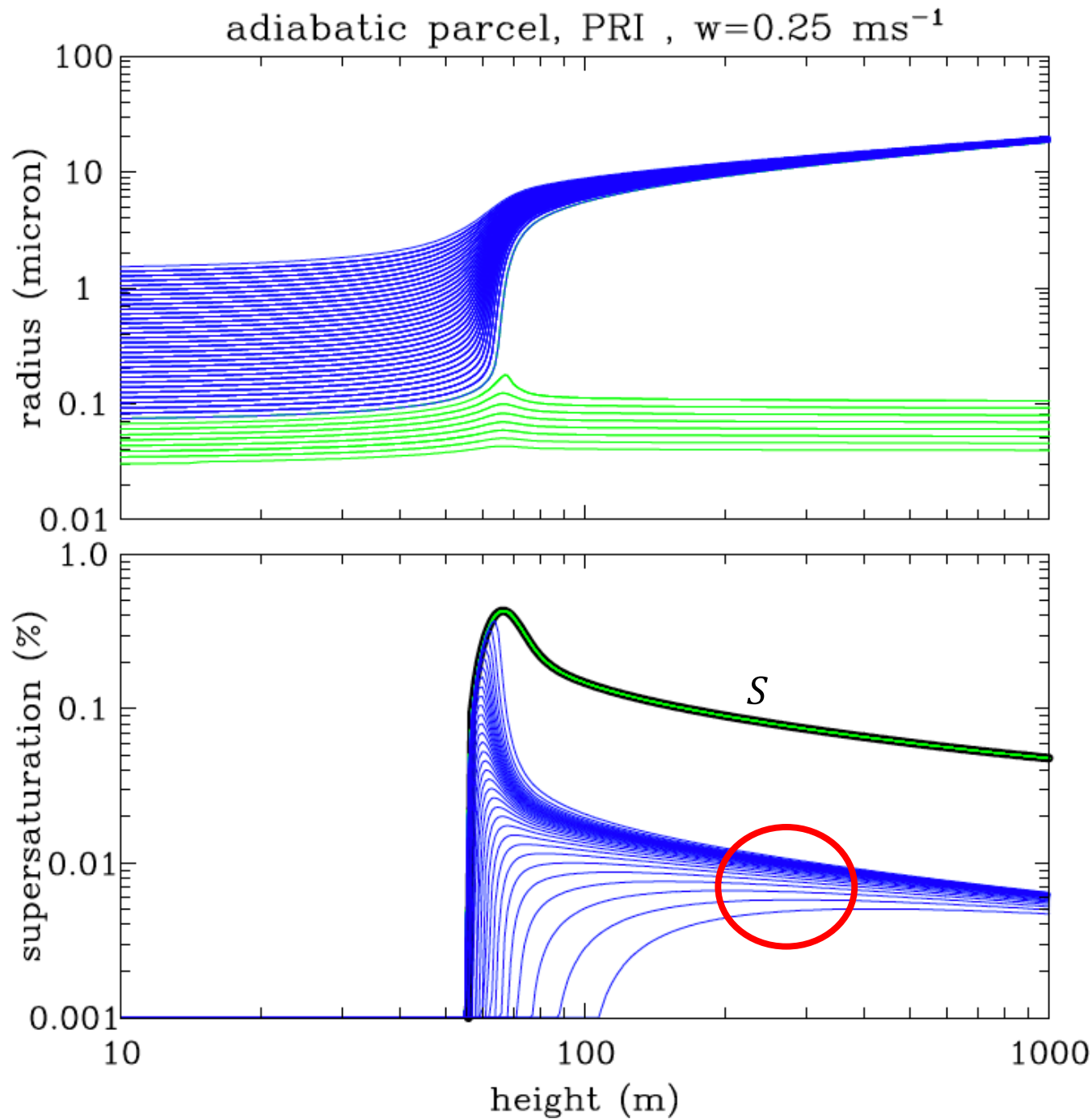


S

$$\frac{dr^2}{dt} \sim (S - S_{eq})$$

spectral width in r^2 space:

$$\begin{aligned} & d(r_1^2 - r_2^2)/dt \\ &= -G [S_{eq}(r_1) - S_{eq}(r_2)] \end{aligned}$$

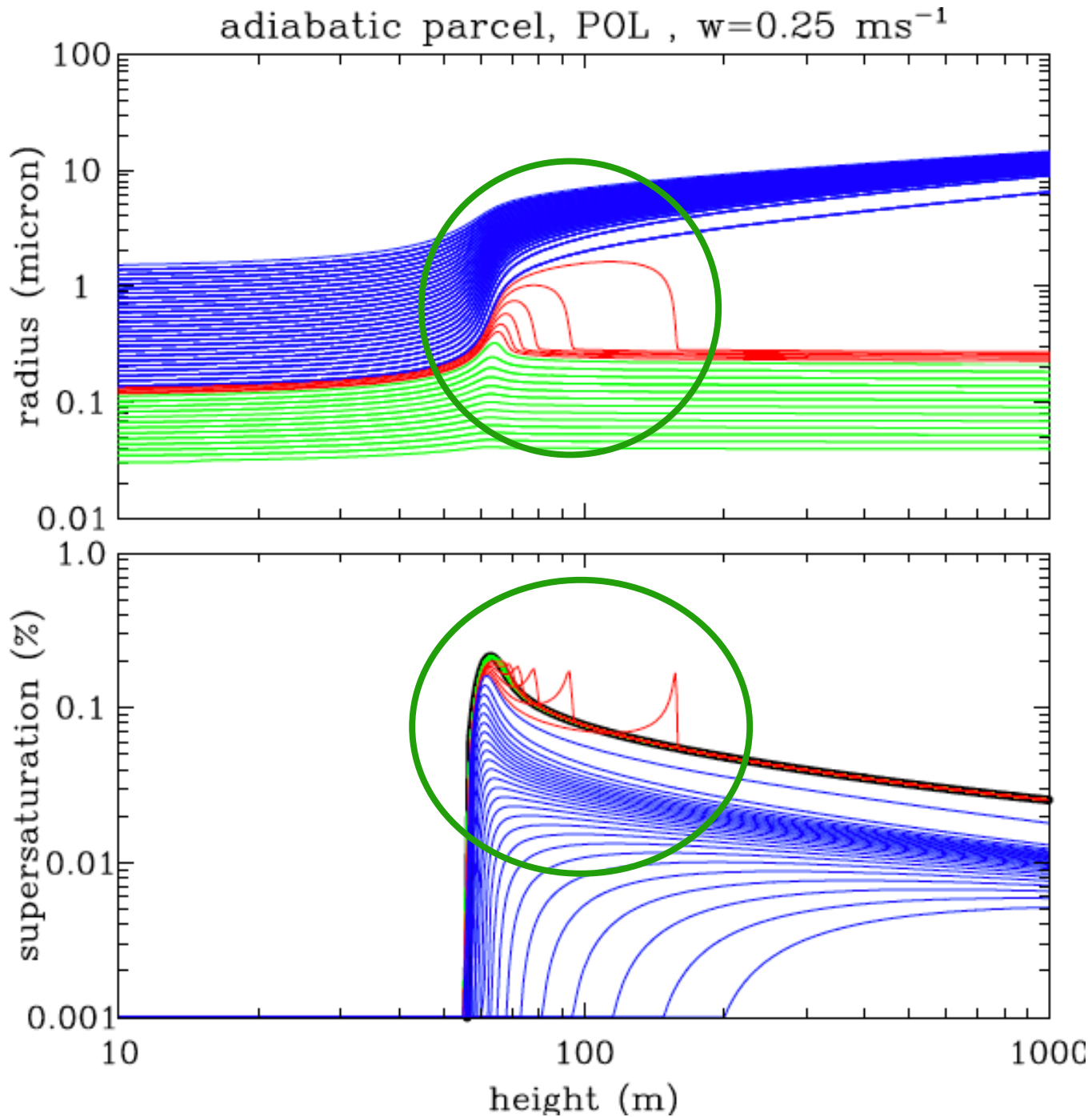


$$\frac{dr^2}{dt} \sim (S - S_{eq})$$

spectral width in r^2 space:

$$\begin{aligned} & d(r_1^2 - r_2^2)/dt \\ &= -G [S_{eq}(r_1) - S_{eq}(r_2)] \end{aligned}$$

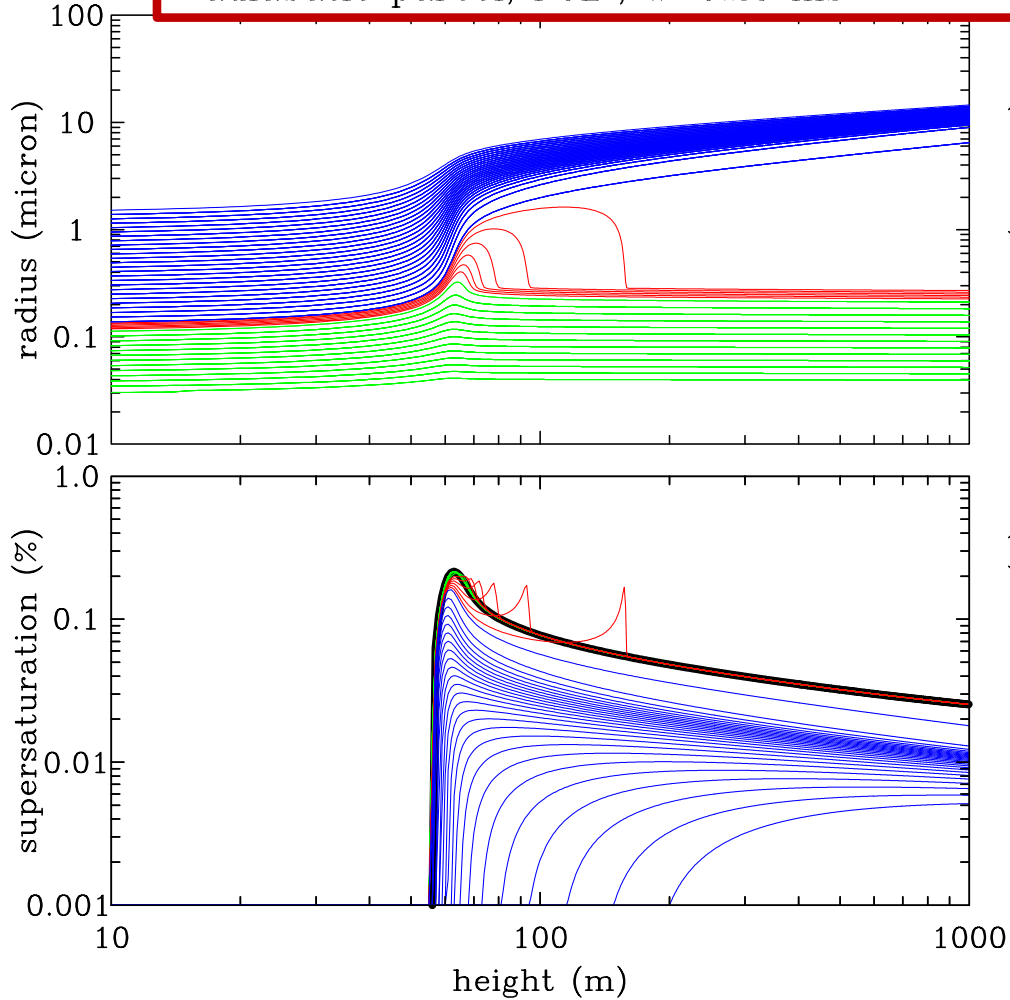
positive if $r_1 > r_2$!!!!



$$\frac{dr^2}{dt} \sim (S - S_{eq})$$

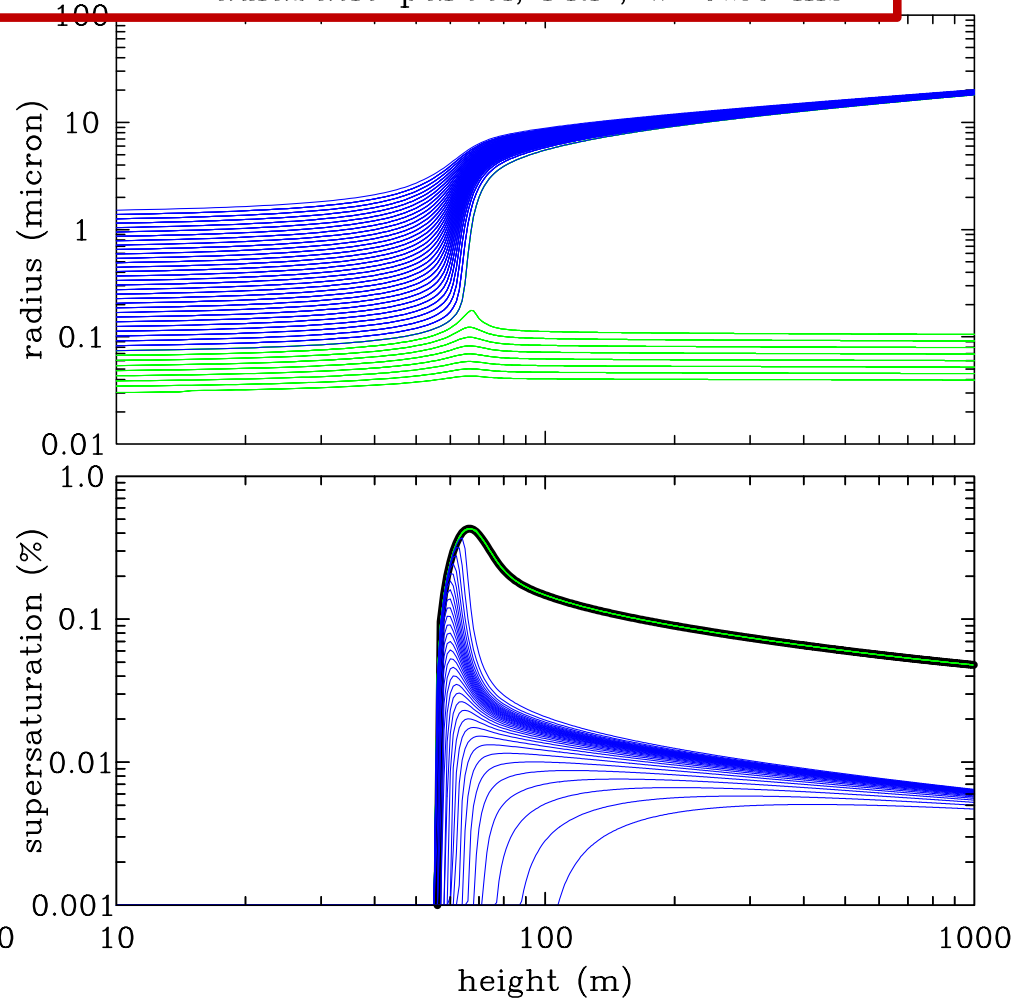
lower maximum
supersaturation in the
polluted case allows
CCN deactivation!

adiabatic parcel, POL , $w=0.25 \text{ ms}^{-1}$



polluted

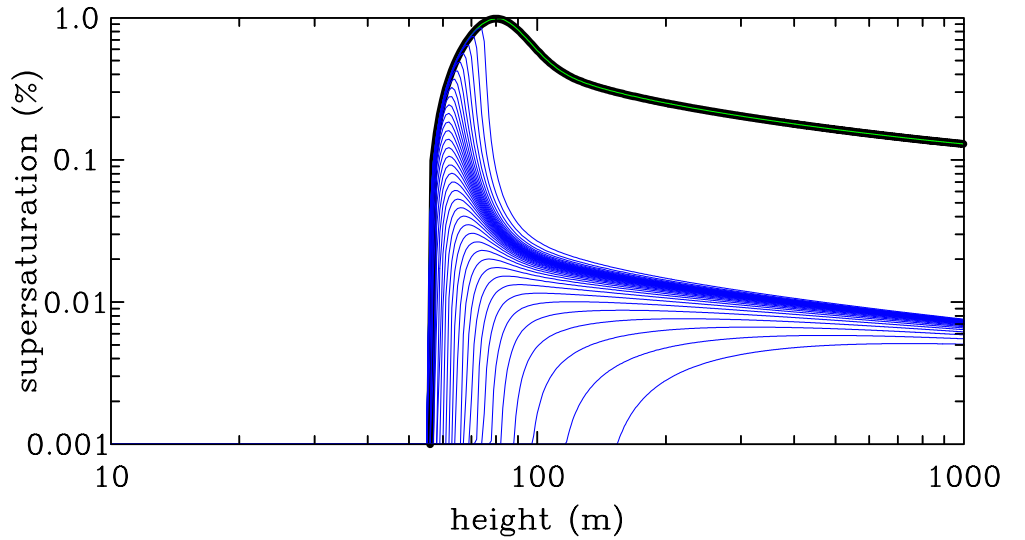
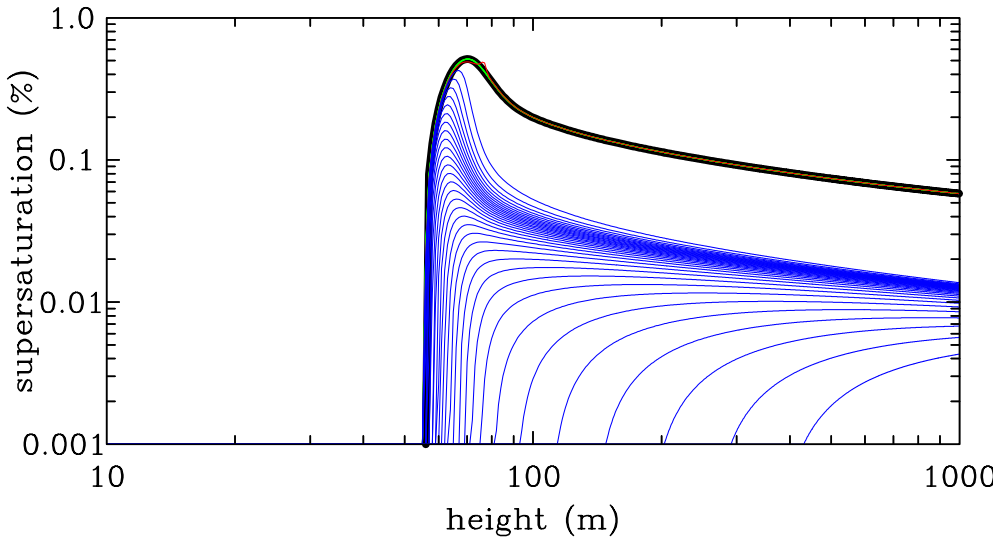
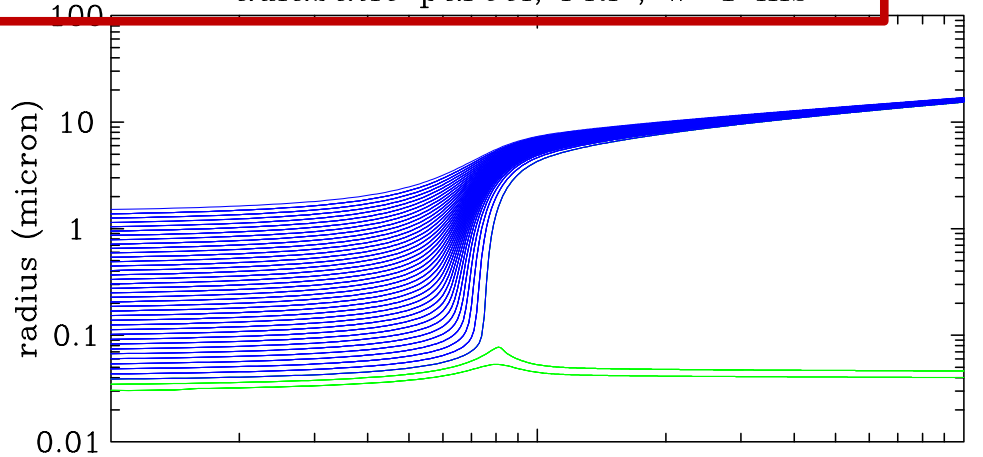
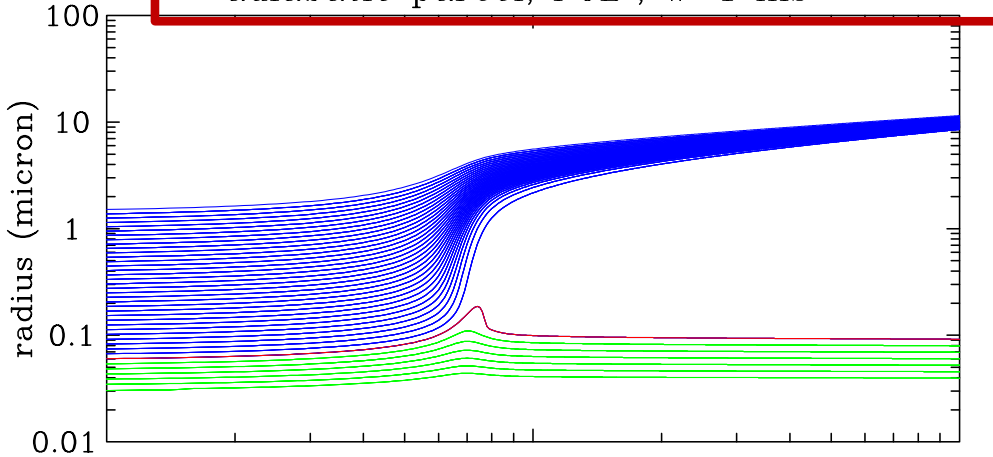
adiabatic parcel, PRI , $w=0.25 \text{ ms}^{-1}$



pristine

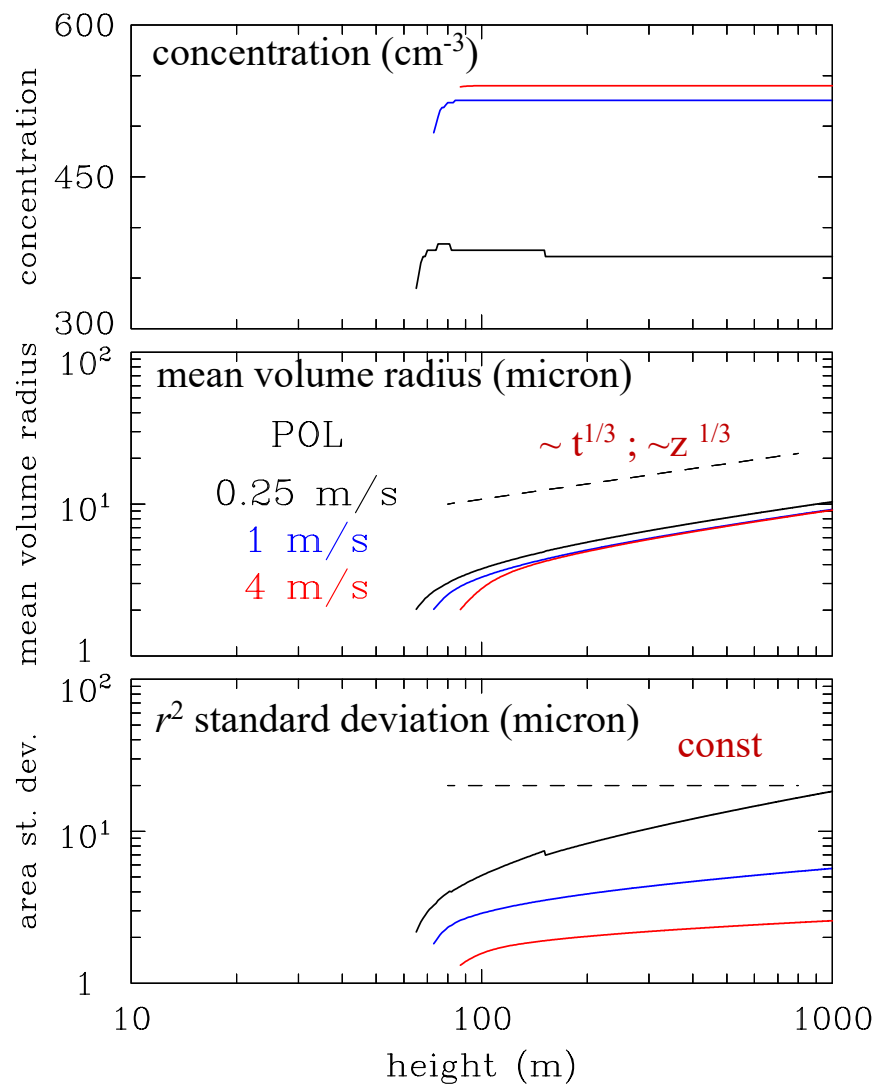
adiabatic parcel, POL , $w=1 \text{ ms}^{-1}$

adiabatic parcel, PRI , $w=1 \text{ ms}^{-1}$

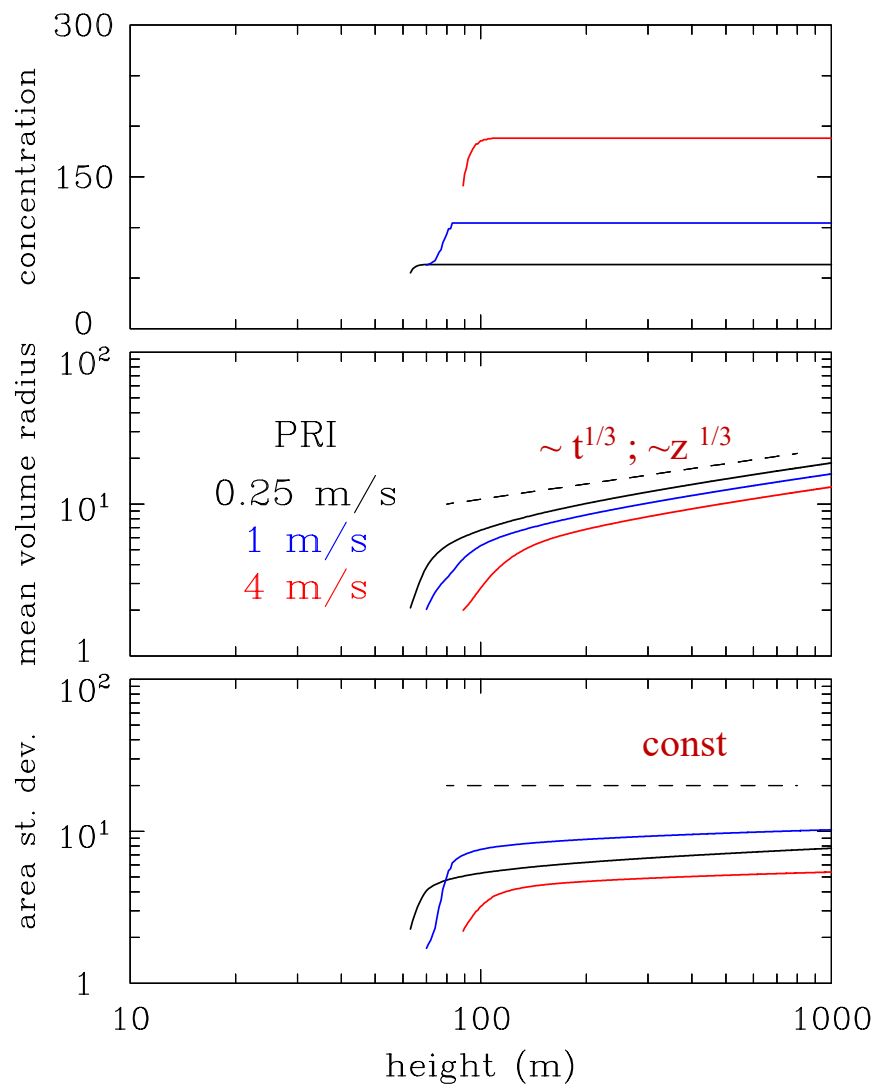


polluted

pristine



polluted



pristine

Radiative properties (optical thickness, etc.) of warm clouds depend on the so-called **effective radius**, the ratio of the third and second moment of the droplet size distribution:

$$R_{eff} = \frac{\langle R^3 \rangle}{\langle R^2 \rangle}$$

Effective radius R_{eff} depends on the **mean volume radius** $\sim(\text{LWC}/N)^{1/3}$ and the droplet spectrum **relative dispersion** d , the ratio of the standard deviation of the droplet radius distribution and the mean droplet radius (Pontikis and Hicks GRL 1992; “PH” below; Liu and Daum GRL 2000):

effective radius

$$R_{eff} = \alpha \left(\frac{L}{N} \right)^{1/3}$$

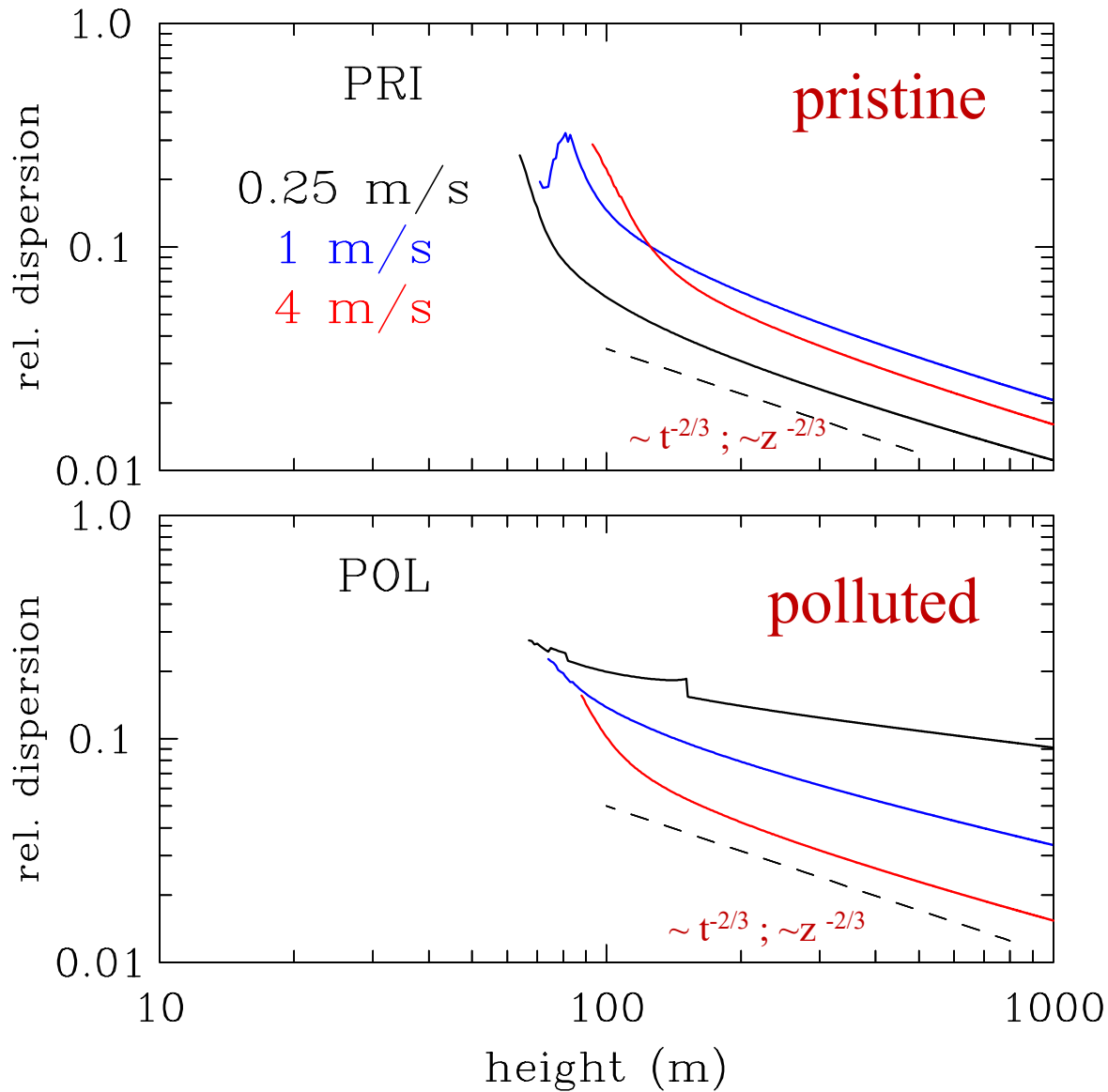
L – LWC

N – droplet concentration

$$d = \sigma / \langle R \rangle$$

relative dispersion

$$\alpha_{PH}(d) = 62.04 \frac{(1 + 3d^2)^{2/3}}{(1 + d^2)}$$



Liu and Daum (GRL 2000):

assuming $dr^2/dt \sim S$, $d \sim r^2$

so $d \sim z^{-2/3}$ because $r \sim z^{1/3}$

$d = \sigma / \langle R \rangle$ - relative dispersion

The role of CCN distribution:

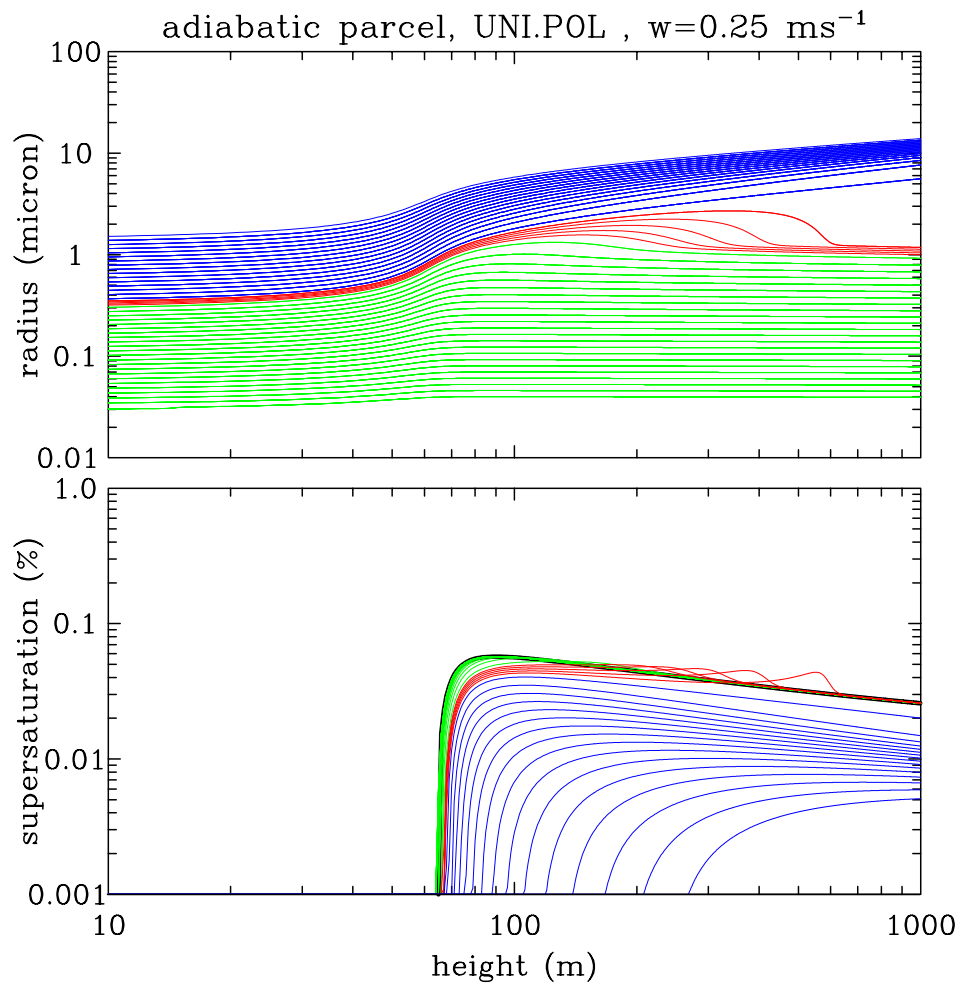
additional simulations with different distributions:

a uniform CCN distribution (applying 200 uniformly distributed classes in the logarithmic space in the range between 4 and 400 nm as in all other simulations):

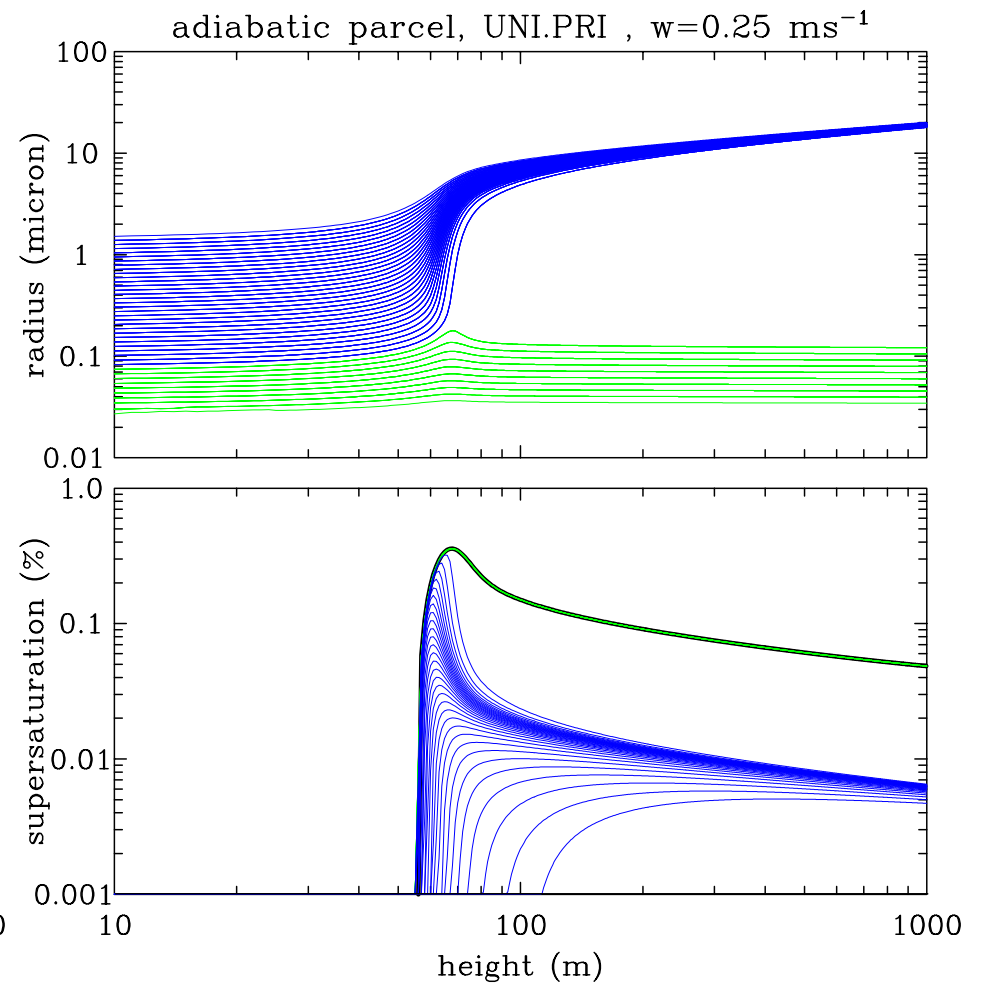
total concentration of 100 cm^{-3} , set UNI.PRI

total concentration of $1,000 \text{ cm}^{-3}$, set UNI.POL

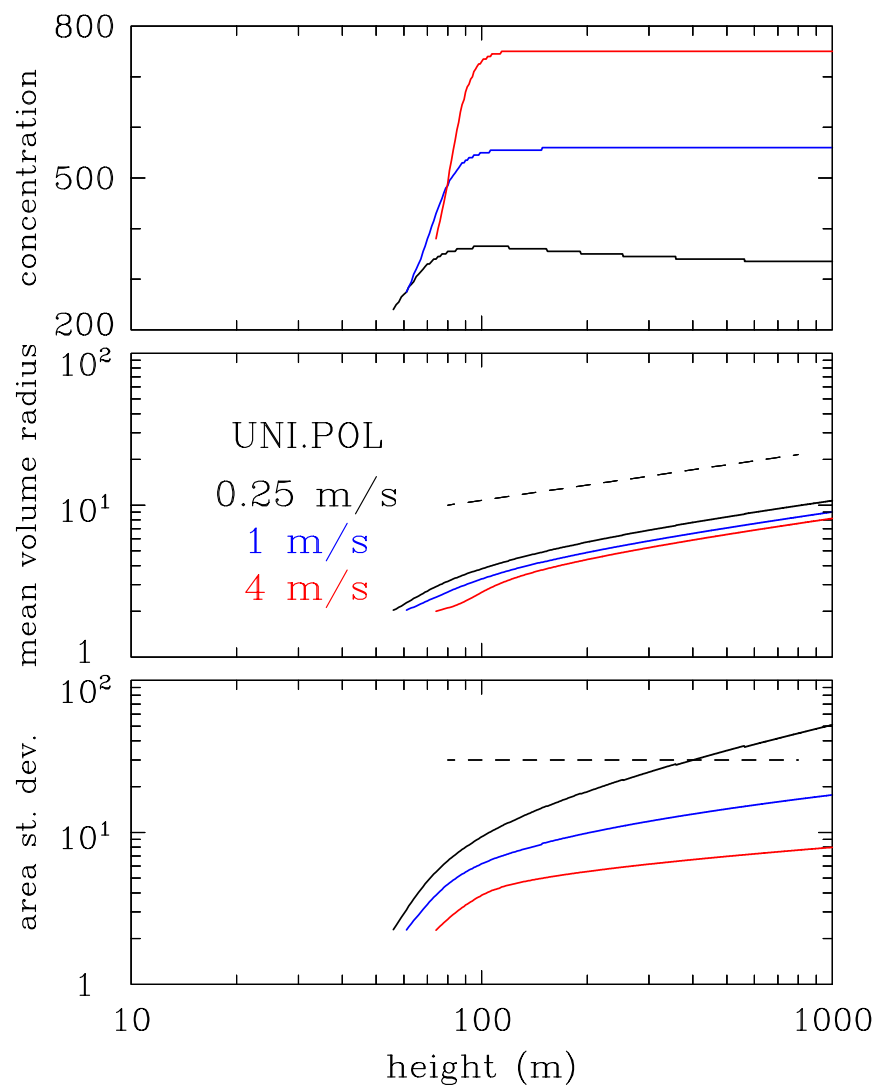
This implies that each dry CCN size class has a concentration of either 0.5 cm^{-3} (UNI.PRI) or 5 cm^{-3} (UNI.POL).



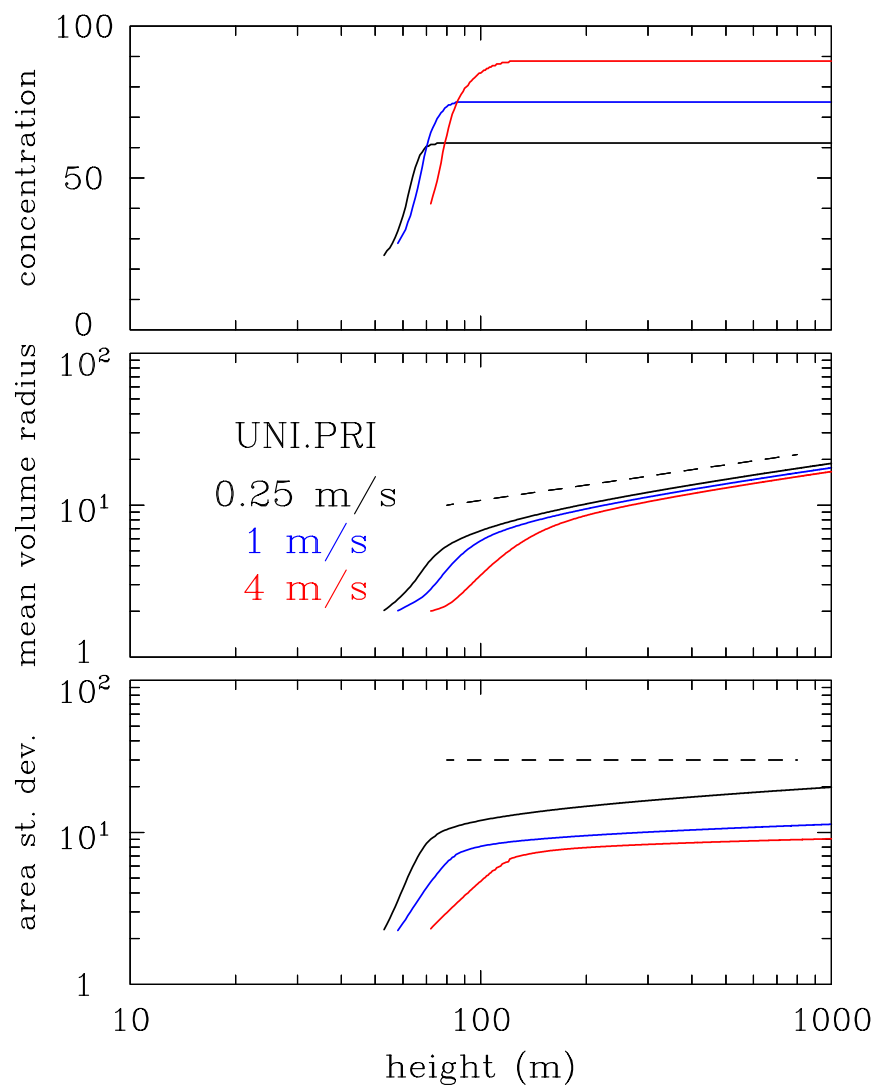
polluted



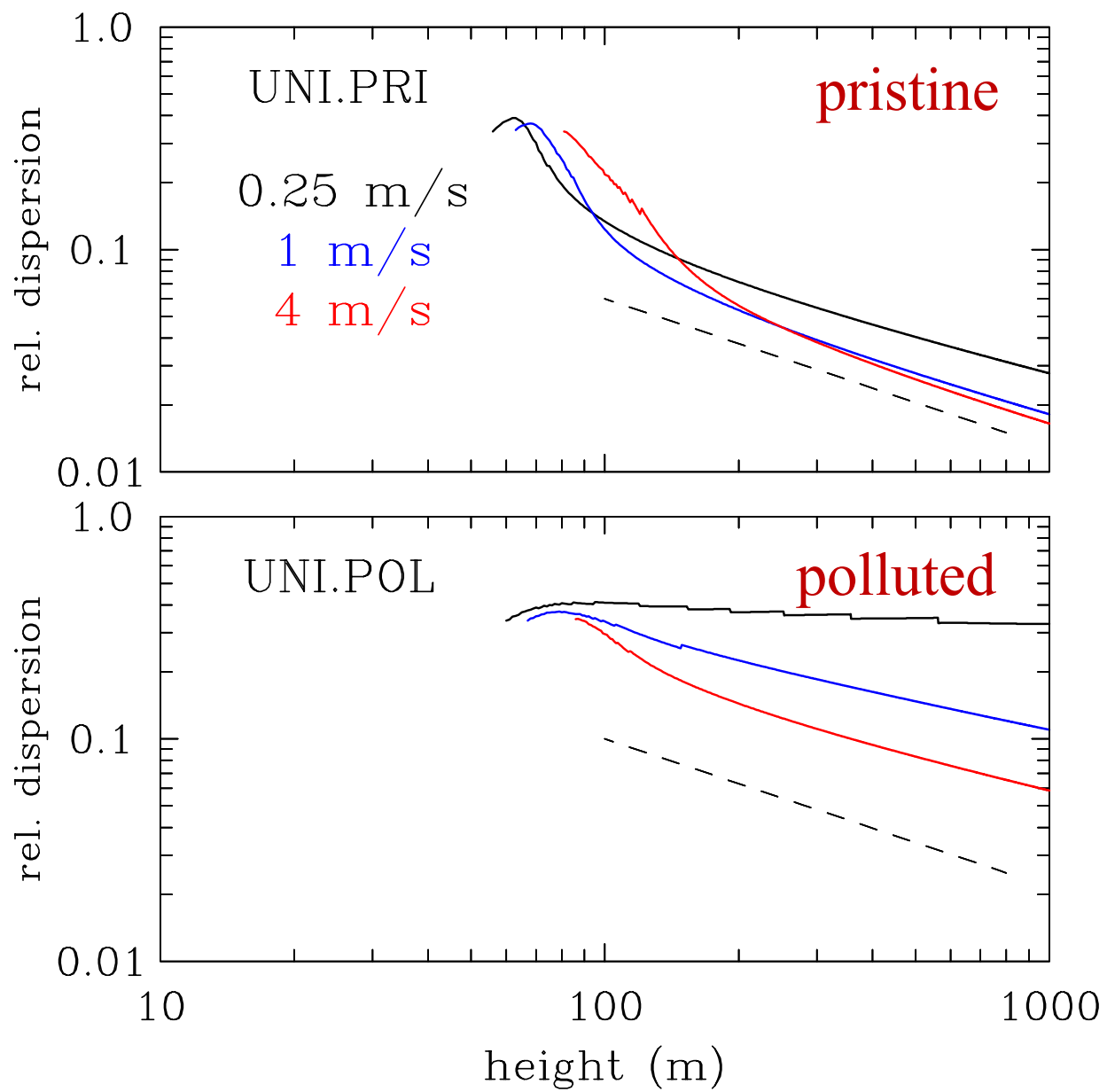
pristine



polluted

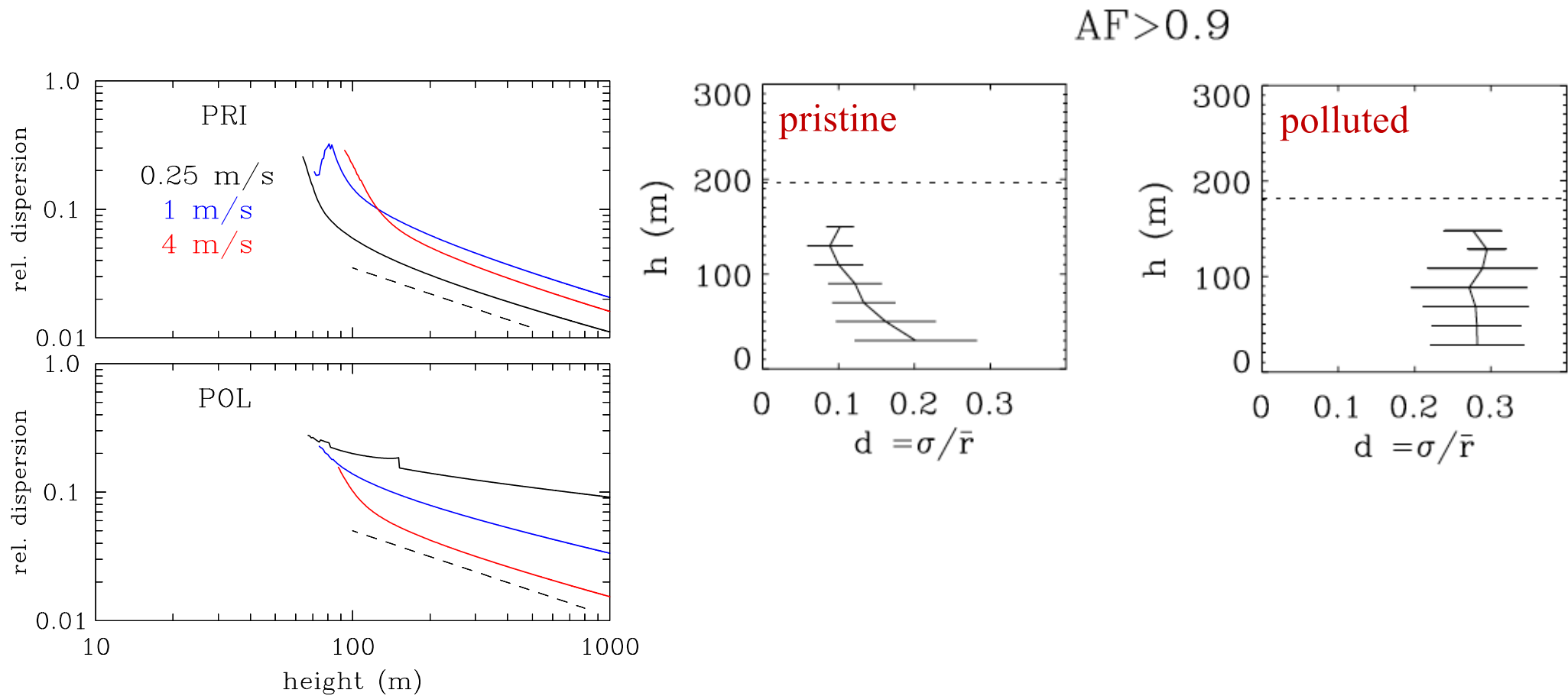


pristine



Observations of the width of cloud droplet spectra in stratocumulus

H. Pawlowska,¹ W. W. Grabowski,² and J.-L. Brenguier³



Summary (key points from the paper):

Evolutions of cloud droplet spectral width in an adiabatic parcel rising through cloud base differ between clean and polluted environments.

Smaller droplet sizes make the solute and surface tension effects more influential for spectral width evolutions in polluted clouds.

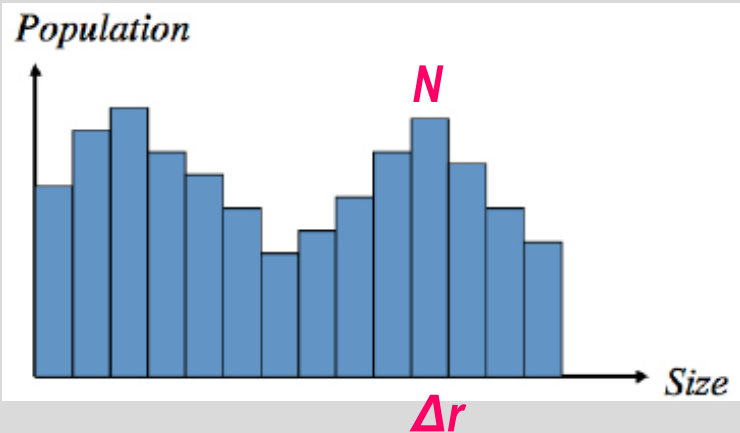
Conventional bin microphysics are not suitable for studying weak-updraft shallow polluted clouds such as continental stratocumulus.

BIN-RESOLVING WARM MICROPHYSICS:

Introducing *spectral density function* $f(r, t)$:

$$f(r, t) \equiv \frac{dN(r, t)}{dr}$$

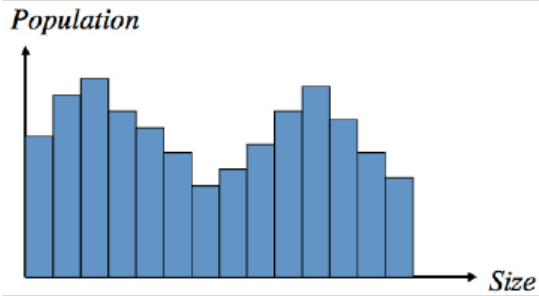
$dN(r, t)$ is the concentration (per unit mass as mixing ratio) of droplets smaller than r (cumulative concentration).



$$f = N / \Delta r$$

BIN-RESOLVING WARM MICROPHYSICS:

ACTIVATION AND CONDENSATION

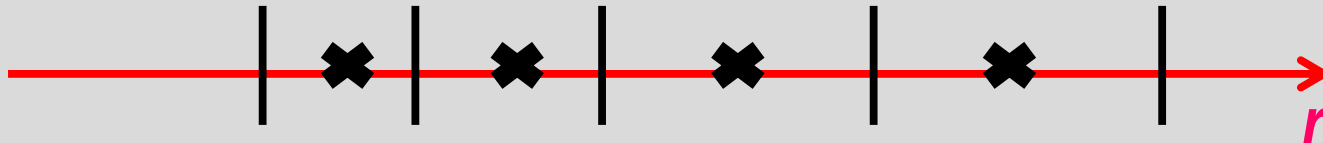


Continuity equation for activation and growth by condensation:

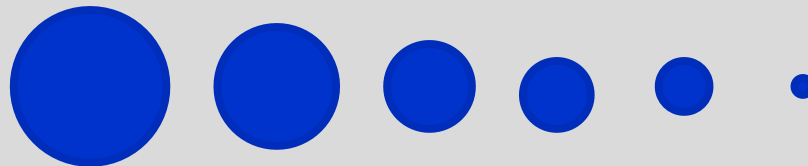
$$\frac{\partial f(r, t)}{\partial t} + \frac{\partial}{\partial r} \left(\frac{dr}{dt} f(r, t) \right) = S_{nucl} \quad \frac{dr}{dt} \sim \frac{S}{r}$$

where S_{nucl} is the source associated with activation of cloud droplets (CCN activation).

cloud droplets

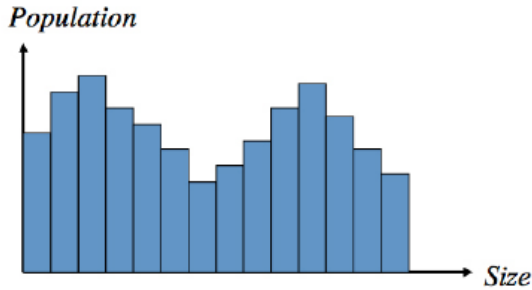


aerosols (CCN)



BIN-RESOLVING WARM MICROPHYSICS:

ACTIVATION AND CONDENSATION

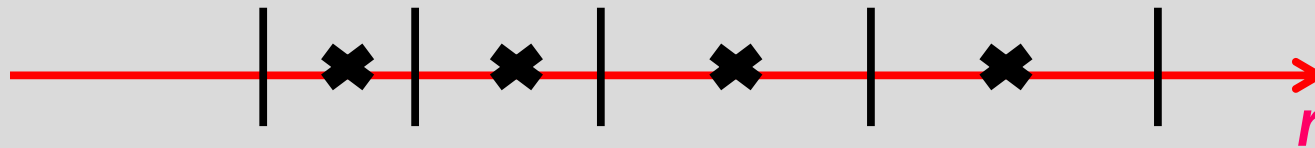


Continuity equation for activation and growth by condensation:

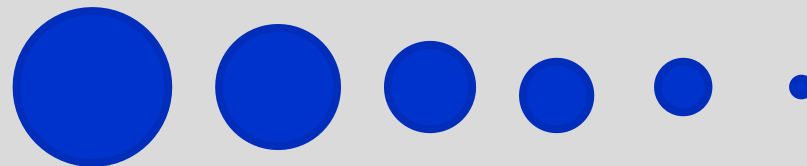
$$\frac{\partial f(r,t)}{\partial t} + \frac{\partial}{\partial r} \left(\frac{dr}{dt} f(r,t) \right) = S_{nucl}$$

where S_{nucl} is the source associated with activation of cloud droplets (CCN activation).

cloud droplets



aerosols (CCN)



move activated
CCN to
droplet grid
once
activated...